

Compositeness of near-threshold s-wave resonances



Tetsuo Hyodo

Tokyo Metropolitan Univ.

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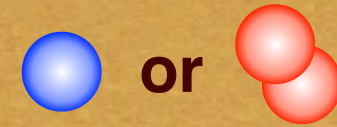
Contents

1 Introduction — threshold rule?

Compositeness

S. Weinberg, Phys. Rev. 137, B672 (1965);

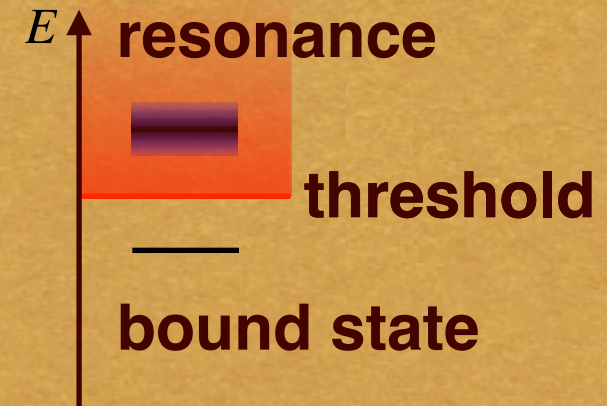
T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 (2013)



Near-threshold bound states

T. Hyodo, PRC90, 055208 (2014);

T. Kinugawa, T. Hyodo, PRC 109, 045205 (2024)



Near-threshold resonances

T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

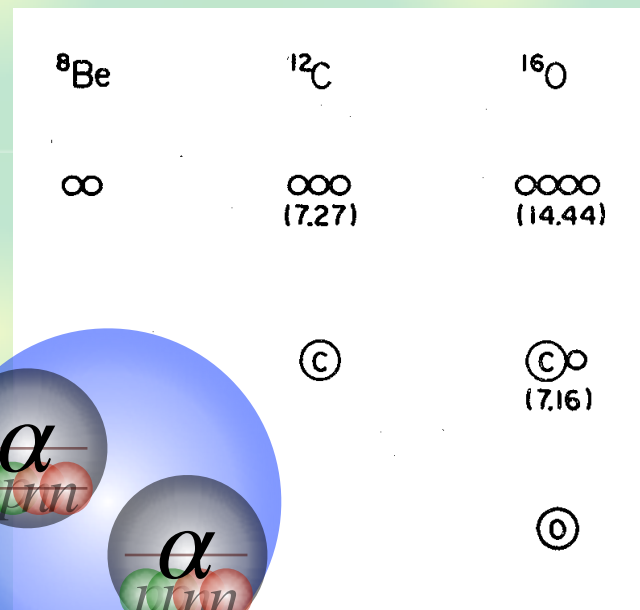
Summary

Molecule-like structure and threshold rule

α cluster in nuclei

H. Horiuchi, K. Ikeda, Y. Suzuki, PTPS 55, 89 (1972)

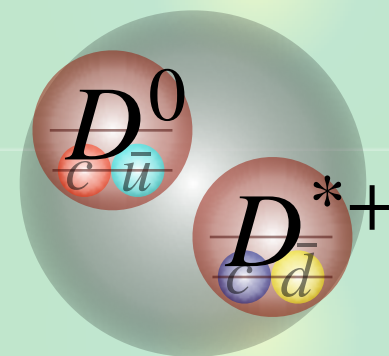
- cluster formation near $n\alpha$ thresholds
- e.g.) ${}^8\text{Be} \sim \alpha\alpha$, ${}^{12}\text{C}$ Hoyle state $\sim \alpha\alpha\alpha$, ...
- molecule-like structure



Hadronic molecules

F.K. Guo, et al., RMP 90, 015004 (2018)

- Exotic hadrons near two-hadron thresholds
- e.g.) $T_{cc} \sim D^0 D^{*+}$, $X(3872) \sim D^0 \bar{D}^{*0}$, ...



Threshold rule:

- molecule-like structure appears near threshold

Questions on threshold rule

Mechanism which leads to threshold rule

- **low-energy universality** ($|a_0| \rightarrow \infty$ for $B \rightarrow 0$)

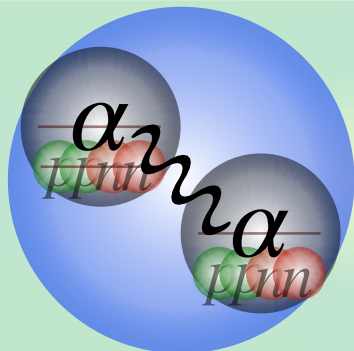
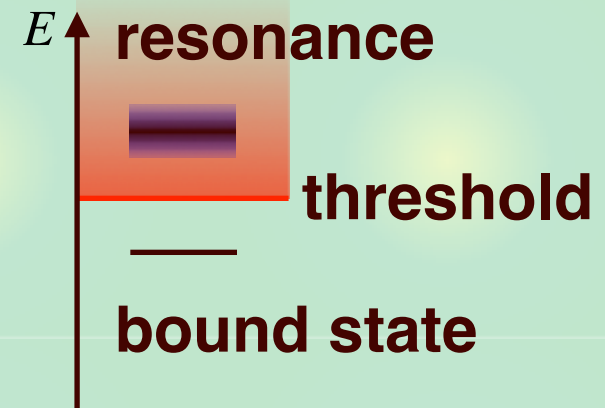
E. Braaten, H.-W. Hammer, *Phys. Rept.* **428**, 259 (2006);

P. Naidon, S. Endo, *Rept. Prog. Phys.* **80**, 056001 (2017)

- **spatially spread wavefunction** \rightarrow molecule

Typical example: ${}^8\text{Be} \sim \alpha\alpha$

- $E_R \sim +0.1$ MeV, **above threshold**
- **wavefunction of unstable resonance?**
- $Z_\alpha = 2$: **Coulomb interaction?**



\rightarrow Next talk by Kinugawa-san

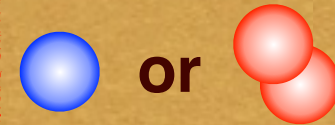
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Near-threshold bound states

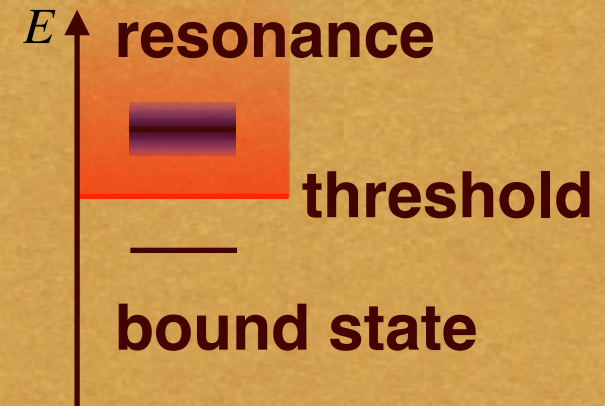
T. Hyodo, PRC90, 055208 (2014);

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Near-threshold resonances

T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

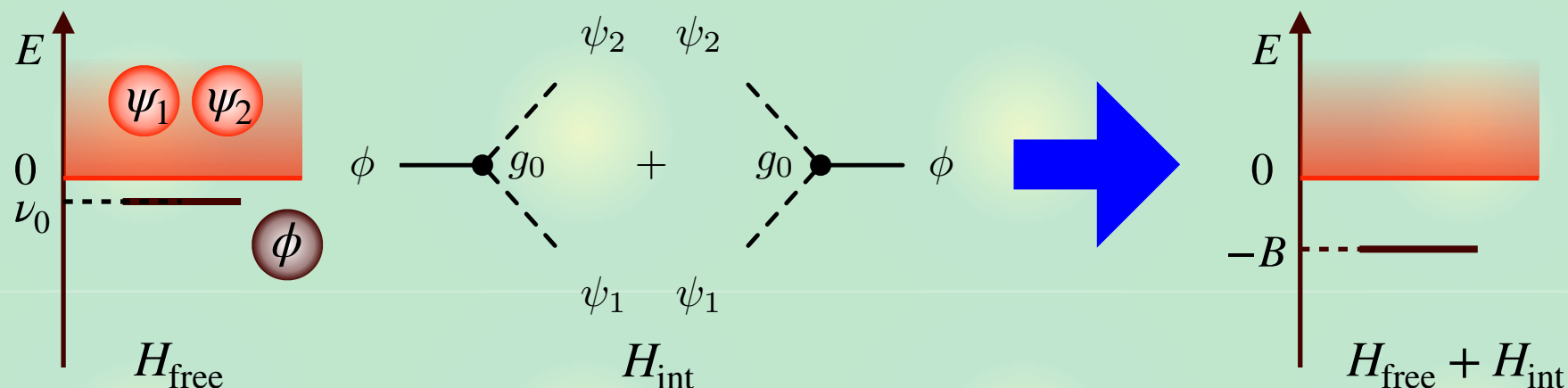
Summary



EFT setup

EFT with bare state ϕ + scattering states $\psi_1\psi_2$

T. Kinugawa, T. Hyodo, PRC 109, 045205 (2024)



- eigenstates of free/full Hamiltonian

$$H_{\text{free}} |B_0\rangle = \nu_0 |B_0\rangle, \quad H_{\text{free}} |\mathbf{p}\rangle = \frac{\mathbf{p}^2}{2\mu} |\mathbf{p}\rangle, \quad (H_{\text{free}} + H_{\text{int}}) |B\rangle = -B |B\rangle$$

- $\psi_1\psi_2$ scattering amplitude: effective range expansion

$$f(p) = -\frac{\mu}{2\pi} \frac{1}{v(E)^{-1} - G(E)} = \frac{1}{-\frac{1}{a_0} + \frac{r_e}{2} p^2 - ip}$$

Compositeness and elementarity

Compositeness: **quantitative** measure of internal structure

- Normalization of $|B\rangle$ + completeness of free eigenstates

$$\langle B|B\rangle = 1, \quad 1 = |B_0\rangle\langle B_0| + \int \frac{d\mathbf{p}}{(2\pi)^3} |\mathbf{p}\rangle\langle \mathbf{p}|$$

- Definition

$$1 = Z + X, \quad Z \equiv |\langle B_0|B\rangle|^2, \quad X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p}|B\rangle|^2$$



“elementarity”

compositeness



- Z, X : real and nonnegative \rightarrow interpreted as **probability**

- Closed-form expressions

$$X = \frac{G'(E)}{[v^{-1}(E)]' - G'(E)} \Bigg|_{E=-B} = 1 - \frac{1}{1 - \Sigma'(E)} \Bigg|_{E=-B}$$

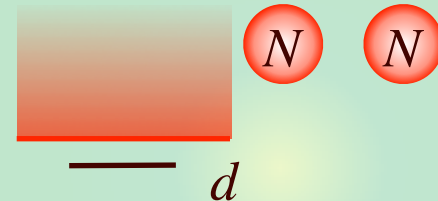
Weak-binding relation

Compositeness X of **stable** bound state

S. Weinberg, *Phys. Rev.* **137**, B672 (1965);

T. Hyodo, *Int. J. Mod. Phys. A* **28**, 1330045 (2013)

$$|d\rangle = \sqrt{X} |NN\rangle + \sqrt{Z} |\text{others}\rangle, \quad X + Z = 1, \quad 0 \leq X \leq 1$$



range of interaction

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad R = \frac{1}{\sqrt{2\mu B}}$$

↑
scattering length

↑
radius of bound state

- $X < 1$ gives violation of universality $a_0 = R$
- for shallow bound state $R \gg R_{\text{typ}}$, $X \leftarrow$ **observables** (a_0, B)

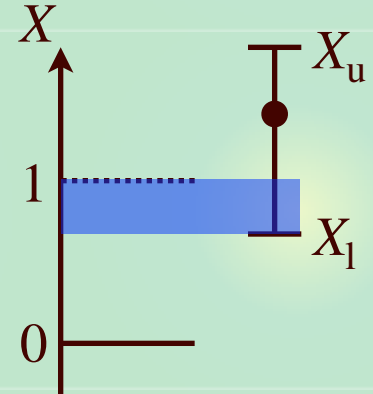
Problem: $a_0 = 5.42$ fm, $R = 4.32$ fm $\Rightarrow X = 1.68 > 1?$

Uncertainty and interpretation

Uncertainty estimation with $\mathcal{O}(R_{\text{typ}}/R)$ term

Y. Kamiya, T. Hyodo, PTEP2017, 023D02 (2017)

$$X_u = \frac{a_0/R + \xi}{2 - a_0/R - \xi}, \quad X_l = \frac{a_0/R - \xi}{2 - a_0/R + \xi}, \quad \xi = \frac{R_{\text{typ}}}{R}$$



Interpretation (with finite range correction)

T. Kinugawa, T. Hyodo, PRC 106, 015205 (2022)

- exclude region outside $0 \leq X \leq 1$

$$R_{\text{typ}} = \max\{R_{\text{int}}, R_{\text{eff}}\}$$

- X of hadrons, **nuclei**, and **atoms**
- X of deuteron is reasonable
- $X \geq 0.5$ in all cases studied

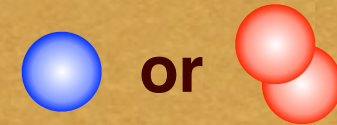
Bound state	Compositeness X
d	$0.74 \leq X \leq 1$
$X(3872)$	$0.53 \leq X \leq 1$
$D_{s0}^*(2317)$	$0.81 \leq X \leq 1$
$D_{s1}(2460)$	$0.55 \leq X \leq 1$
$N\Omega$ dibaryon	$0.80 \leq X \leq 1$
$\Omega\Omega$ dibaryon	$0.79 \leq X \leq 1$
${}^3_{\Lambda}\text{H}$	$0.74 \leq X \leq 1$
${}^4\text{He}$ dimer	$0.93 \leq X \leq 1$

Near-threshold bound states are **mostly composite**

Contents

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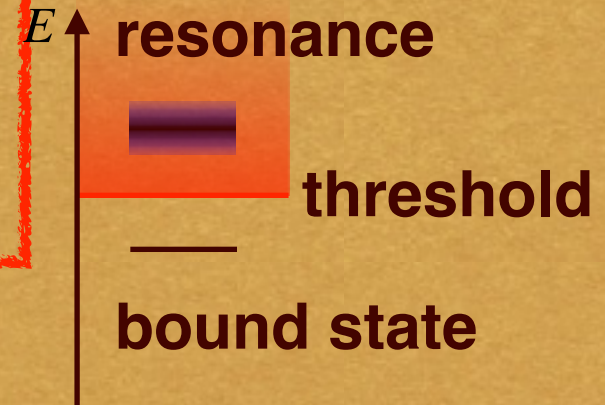
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T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

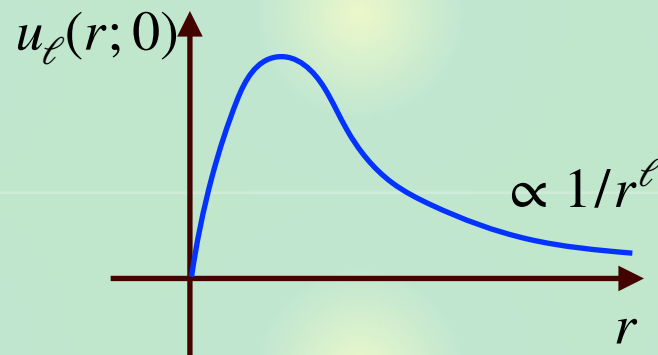
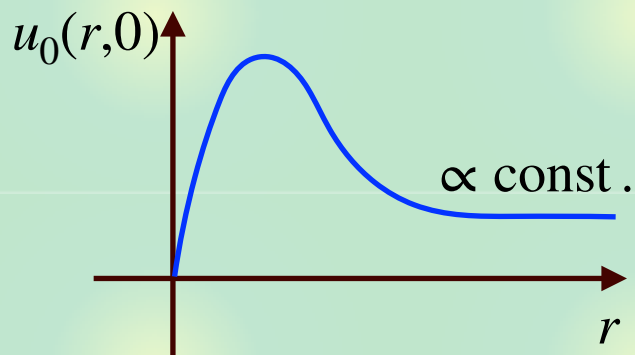
 Summary

compositeness theorem and intuitive picture

Compositeness theorem: $X = 1$ in $B \rightarrow 0$ limit

T. Hyodo, PRC90, 055208 (2014)

- **wavefunction of $B = 0$ state is not normalizable ($|a_0| \rightarrow \infty$)**



- **compositeness in coordinate space**

$$1 = |\langle B_0 | B \rangle|^2 + \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p} | B \rangle|^2 = |\langle B_0 | B \rangle|^2 + \int d\mathbf{r} |\Psi(\mathbf{r})|^2$$

$$\frac{|\langle B_0 | B \rangle|^2}{\int d\mathbf{r} |\Psi(\mathbf{r})|^2} \rightarrow 0 \quad \Rightarrow \quad X = 1$$

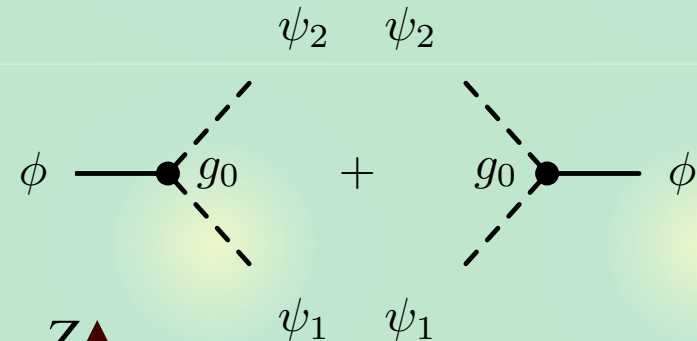
$\rightarrow \infty$

Finite binding case

Elementarity of bound state with small but finite B

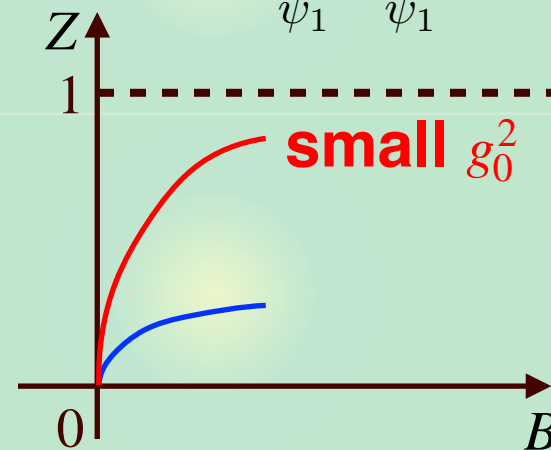
$$Z = \frac{1}{1 - \Sigma'(-B)}$$

$$\sim \frac{1}{1 + Cg_0^2/\sqrt{B}} \sim \frac{\sqrt{B}}{Cg_0^2} + \dots \neq 0$$



B dependence of Z

- $\lim_{B \rightarrow 0} Z = 0$ is fixed
- $Z \ll 1$ for small B (composite)



For sufficiently **small** g_0^2 , $\sqrt{B}/g_0^2 \sim \mathcal{O}(1)$ for small B

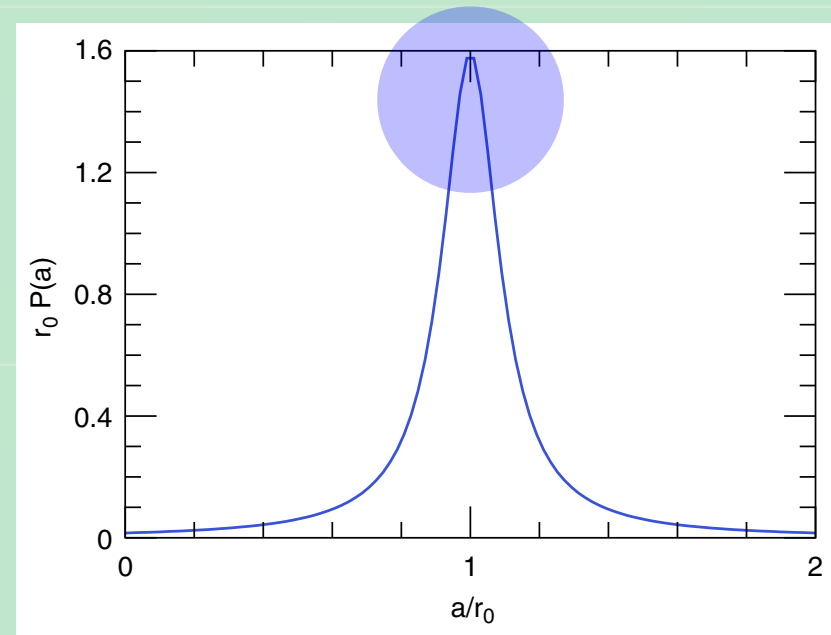
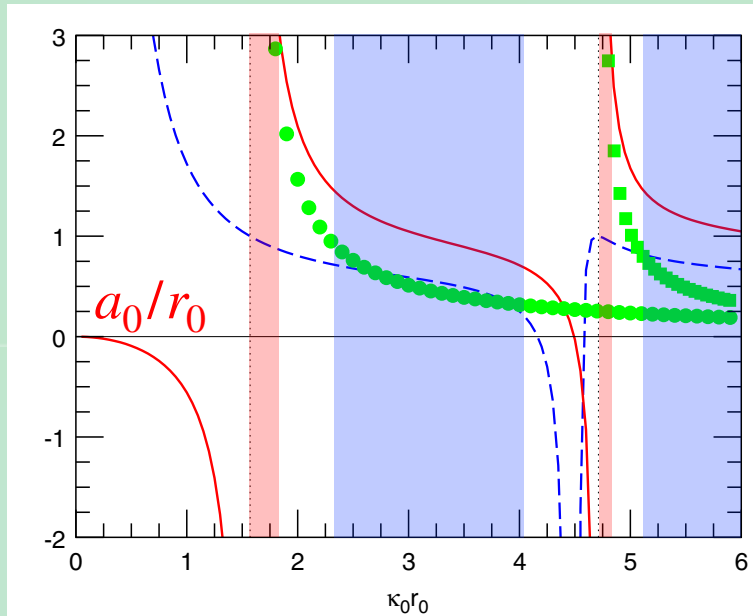
—> sizable Z with small B by **fine tuning** of parameter g_0^2

How probable is such fine tuning?

Quantifying fine tuning

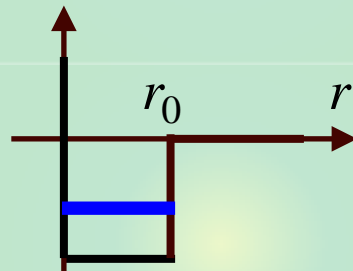
Probability distribution of a_0 of square-well potential

E. Braaten, H.-W. Hammer, Phys. Rept. 428, 259 (2006)



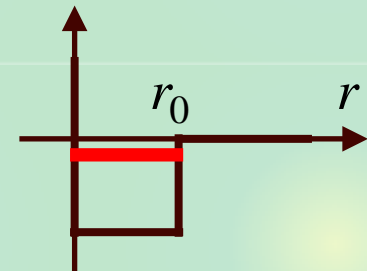
typical

$a_0/r_0 \sim 1$



shallow

$a_0/r_0 \gg 1$

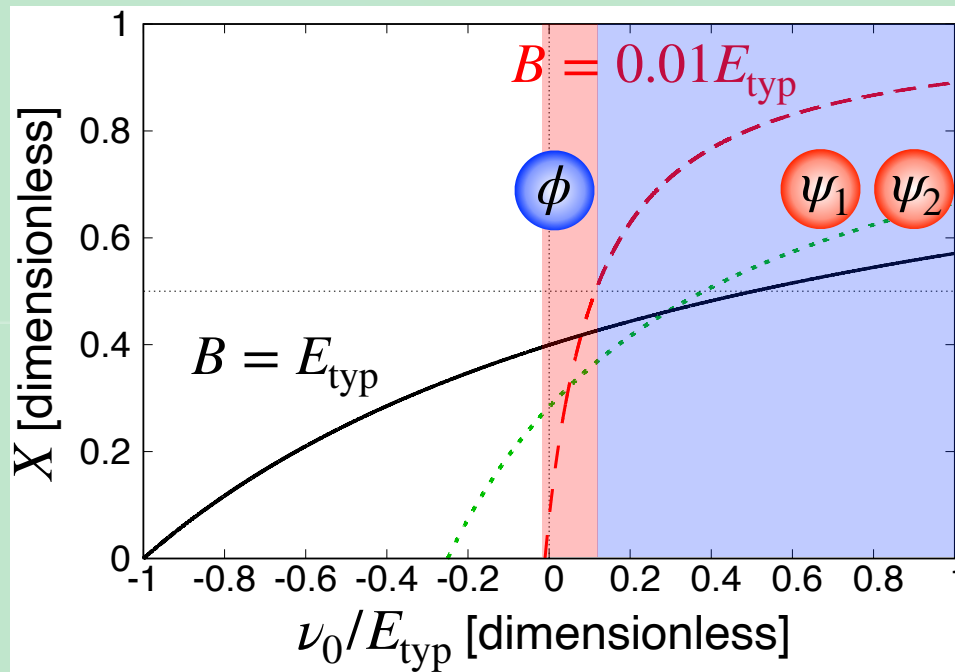


Fine-tuning can be quantified by **parameter dependence**

Structure of bound state

Compositeness X in the allowed ν_0 region

T. Kinugawa, T. Hyodo, PRC 109, 045205 (2024)



- Typical bound state $B = E_{\text{typ}}$: mostly elementary
- Shallow bound state $B = 0.01 E_{\text{typ}}$: mostly **composite**

Shallow elementary state \leftarrow only with **fine tuning** = unlikely

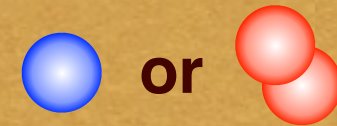
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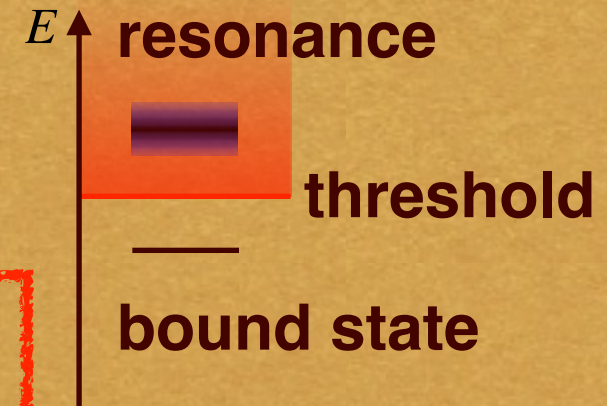
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 Near-threshold resonances

T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

 Summary

Resonances in effective range expansion

Two poles in effective range expansion (ERE)

E. Braaten, H.-W. Hammer, Phys. Rept. 428, 259 (2006);

T. Hyodo, PRL111, 132002 (2013);

T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

$$k^\pm = \frac{i}{r_e} \pm \frac{1}{r_e} \sqrt{\frac{2r_e}{a_0} - 1 + i0^+}$$

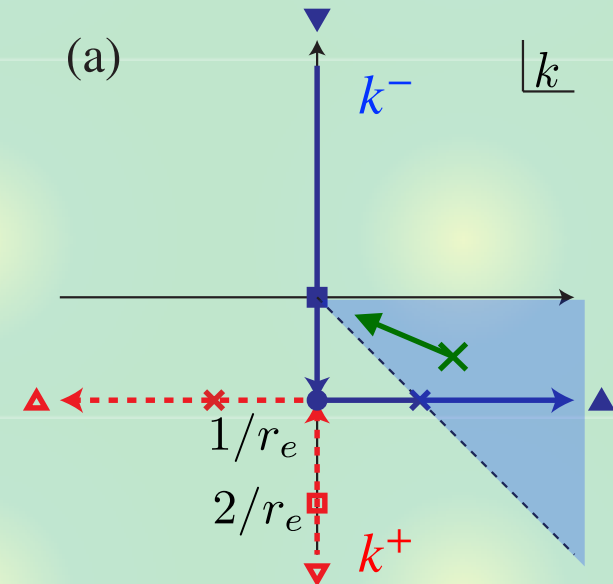
- pole positions $k^\pm \longleftrightarrow (a_0, r_e)$

Resonance solution ($r_e < 0$)

$$\frac{1}{|r_e|} \sqrt{\frac{2r_e}{a_0} - 1} \geq \frac{1}{|r_e|}, \quad \Rightarrow \quad \frac{r_e}{a_0} \geq 1, \quad \Rightarrow \quad |a_0| \leq |r_e|$$

- resonance with $|k^-| \rightarrow 0$: not only $|a_0| \rightarrow \infty$ but also $|r_e| \rightarrow \infty$

- energy $E_R = M_R - i\frac{\Gamma_R}{2} \longleftrightarrow (a_0, r_e)$



Compositeness of resonances

Compositeness: **pure imaginary** ← weak-binding relation

$$X = \sqrt{\frac{1}{1 - \frac{2r_e}{a_0}}} = -i \tan(\theta_k), \quad k^- = |k^-| e^{i\theta_k}$$

Resonance state: complex eigenenergy

$$H|R\rangle = E_R|R\rangle, \quad E_R = M_R - i\frac{\Gamma_R}{2} \in \mathbb{C}$$

- normalization by **Gamow vector**

$$\langle R|H = \langle R|E_R^*, \quad \langle \tilde{R}|H = \langle \tilde{R}|E_R$$

$$\langle R|R\rangle \rightarrow \infty, \quad \langle \tilde{R}|R\rangle = 1$$

- complex Z and X : probability?

$$Z \equiv \langle \tilde{R}|B_0\rangle\langle B_0|R\rangle \in \mathbb{C}, \quad X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} \langle \tilde{R}|\mathbf{p}\rangle\langle \mathbf{p}|R\rangle \in \mathbb{C}$$

New interpretation scheme

Complex matrix element \leftrightarrow **uncertain** nature of resonances

T. Berggren, PLB33, 547 (1970)

Introduce three probabilities $\mathcal{X} + \mathcal{Y} + \mathcal{Z} = 1$

T. Kinugawa, T. Hyodo, arXiv:2403.12635 [hep-ph]

$$\mathcal{X} = \frac{(\alpha - 1)|X| - \alpha|Z| + \alpha}{2\alpha - 1} \quad \text{certainly finding composite}$$

$$\mathcal{Y} = \frac{|X| + |Z| - 1}{2\alpha - 1} \quad \text{uncertain}$$

$$\mathcal{Z} = \frac{(\alpha - 1)|Z| - \alpha|X| + \alpha}{2\alpha - 1} \quad \text{certainly finding elementary}$$

- **Choice of parameter α : exclude large width resonance**

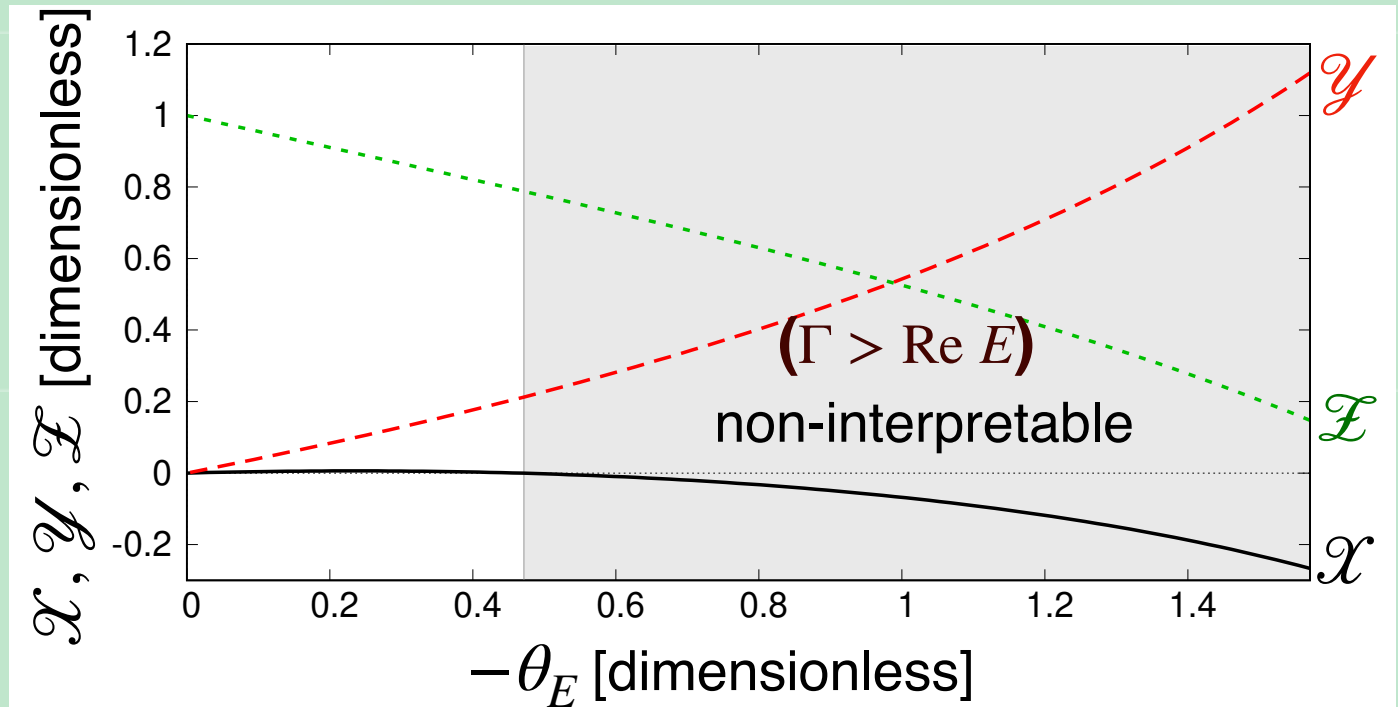
$$\alpha = \frac{\sqrt{5} - 1 + \sqrt{10 - 4\sqrt{5}}}{2} \approx 1.1318$$

- **If $\Gamma > \text{Re } E \rightarrow \mathcal{X} < 0$: non-interpretable state**

Compositeness of resonances

$\mathcal{X}, \mathcal{Y}, \mathcal{Z}$ as functions of argument of eigenenergy

$$E_R = |E_R| e^{i\theta_E}$$






- Resonances are **not composite dominant** ($\mathcal{Z} \gtrsim 0.8$)

← consistent with large negative r_e

Near-threshold resonances are **not composite dominant**


Summary

 **Compositeness X : probability of finding** 

 **Bound state exactly at threshold**


[T. Hyodo, PRC90, 055208 \(2014\);](#)

- **completely composite** $X = 1$

 **Near threshold bound states**

[T. Kinugawa, T. Hyodo, PRC 109, 045205 \(2024\)](#)

- in general, **composite** $X \sim 1$

 **Near-threshold resonances**

[T. Kinugawa, T. Hyodo, arXiv:2403.12635 \[hep-ph\]](#)

- **non-composite**, $\mathcal{X} \lesssim 0.2$

