

Closed-channel parameters of Feshbach resonances

Pascal Naidon, RIKEN



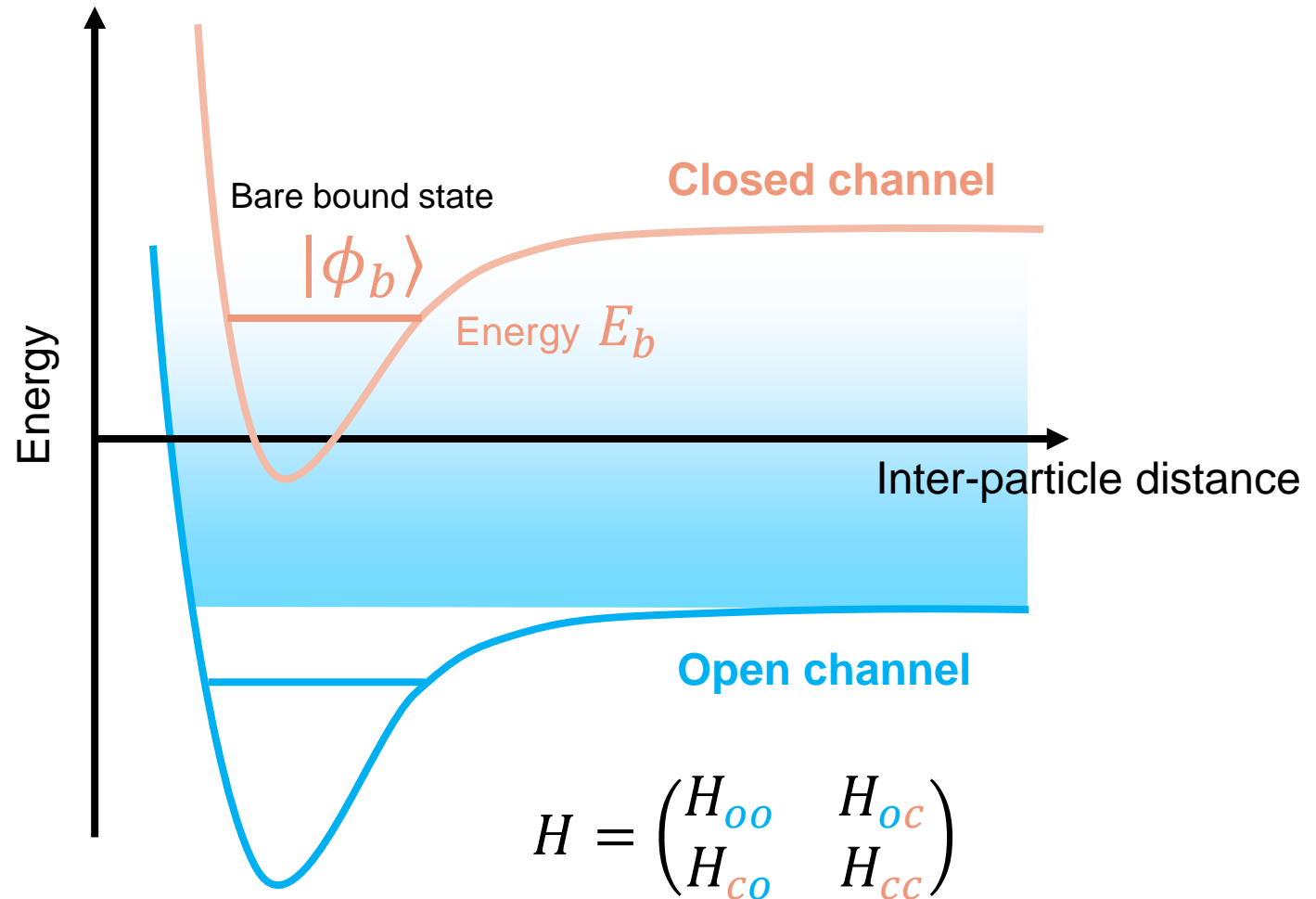
Sep 4, 2024, **Universality of Quantum Systems: From Cold Atoms, Nuclei, to Hadron Physics**
Tohoku University

Motivation

- **Feshbach resonances**: resonance created by a bound state coupled to a continuum.
 - **Ultracold atoms**: very successful tool to control the interactions between ultracold atoms for more than 20 years.
 - **Hadron physics**: hadron resonances near thresholds
- We know how to make models of Feshbach resonances, but **what is physical** and **what is arbitrary** in these models?
- Can we learn about the bound state causing the resonance?

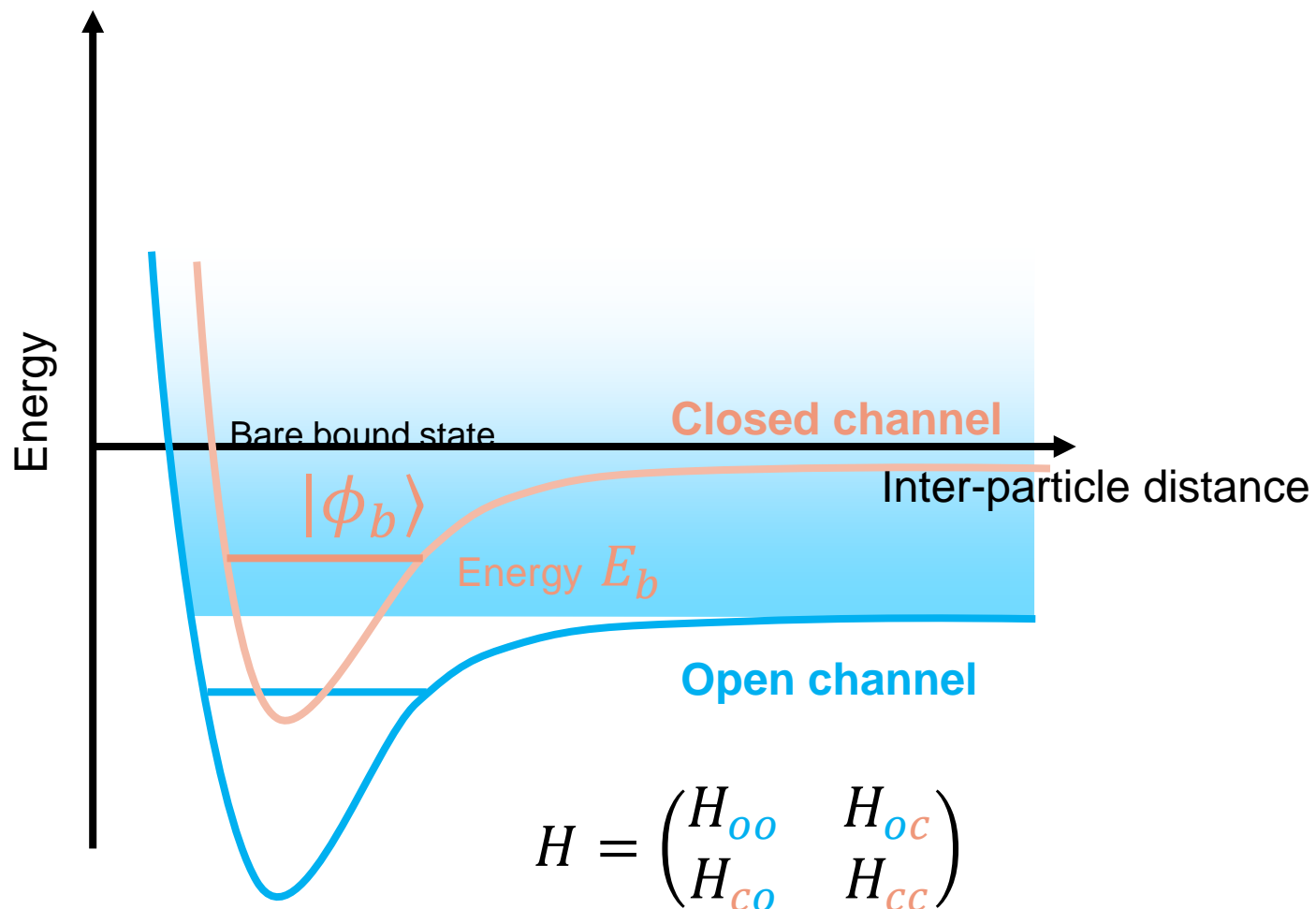
Fano-Feshbach resonances

Principle: two particles can be in different internal states (channels) which are coupled to each other.



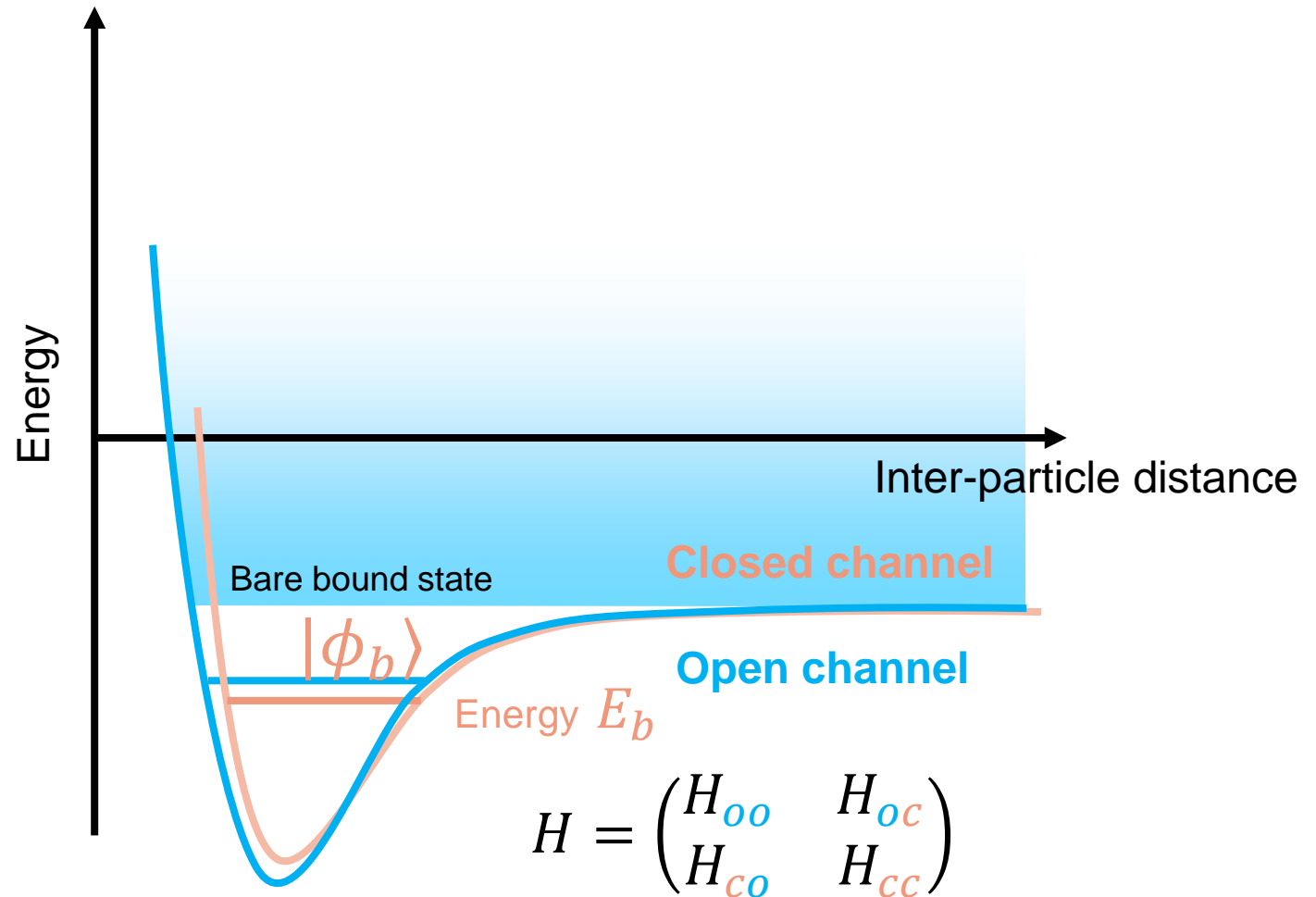
Fano-Feshbach resonances

Principle: two atoms can be in different electronic channels which are coupled to each other.

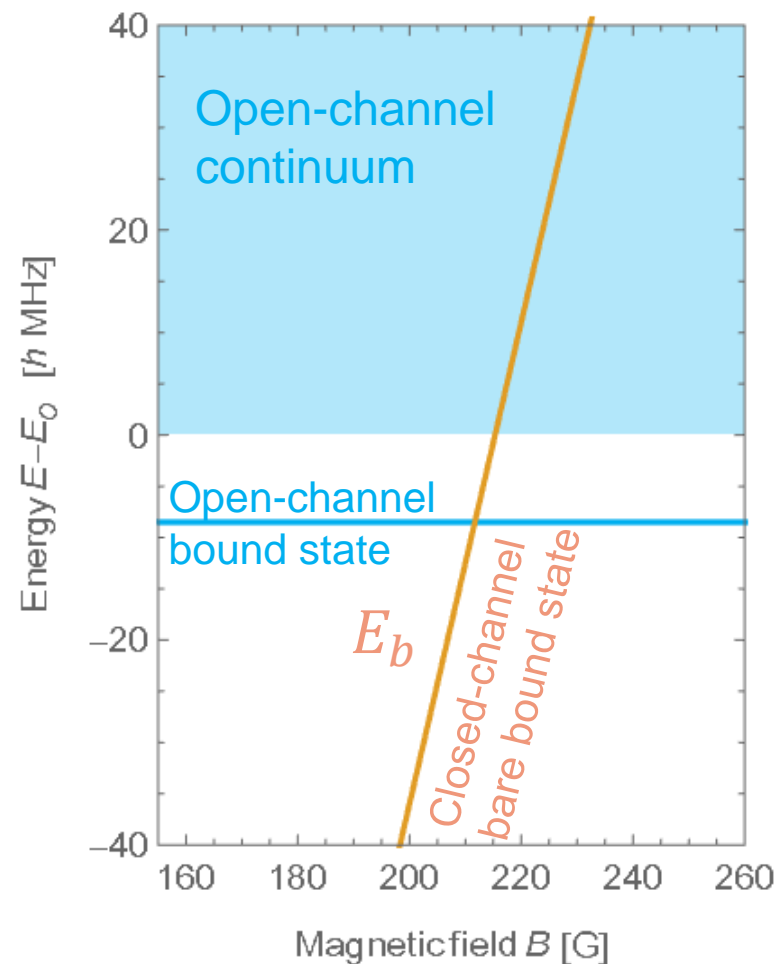


Fano-Feshbach resonances

Principle: two atoms can be in different electronic channels which are coupled to each other.



Two-body observables



Resonance of ^{40}K atoms near 202 G

Two-body observables

$$\gamma = \lim_{k \rightarrow 0} \frac{\Gamma}{2k}$$

$$\Delta_0 = \lim_{k \rightarrow 0} \Delta$$

1) Scattering phase shift:

$$\eta = \eta_0 - \arctan \frac{\Gamma(E)/2}{E - E_b - \Delta(E)}$$

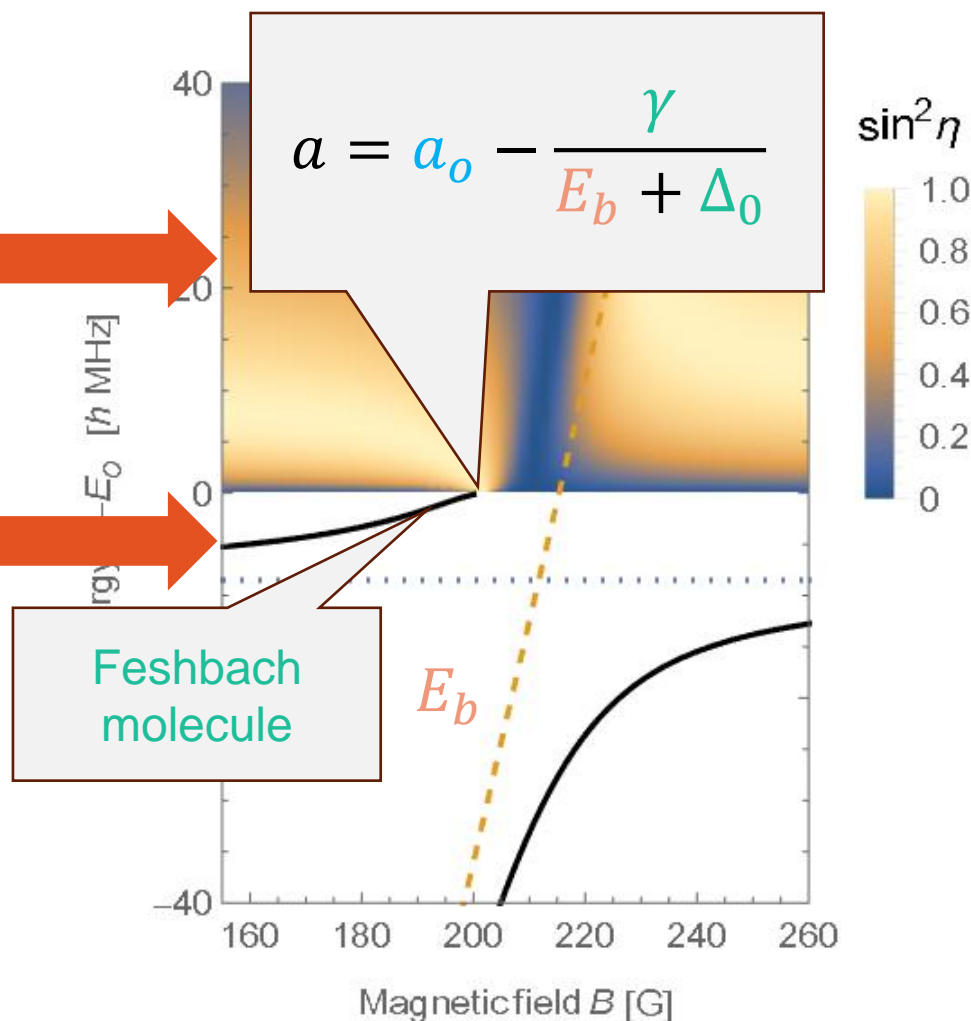
2) Dressed bound-state energy:

$$E = E_b + \Delta(E)$$

Resonance shift Δ and width Γ :

$$\Delta - i\Gamma/2 = \langle \phi_b | H_{co} \frac{1}{E - H_{oo}} H_{oc} | \phi_b \rangle$$

Should depend on the **closed-channel** details!



Resonance of ^{40}K atoms near 202 G

Important question

From these observables, can we determine the **closed-channel** details?

- Energy E_b of the bare bound state ?
- Scattering length a_c in the closed channel ?

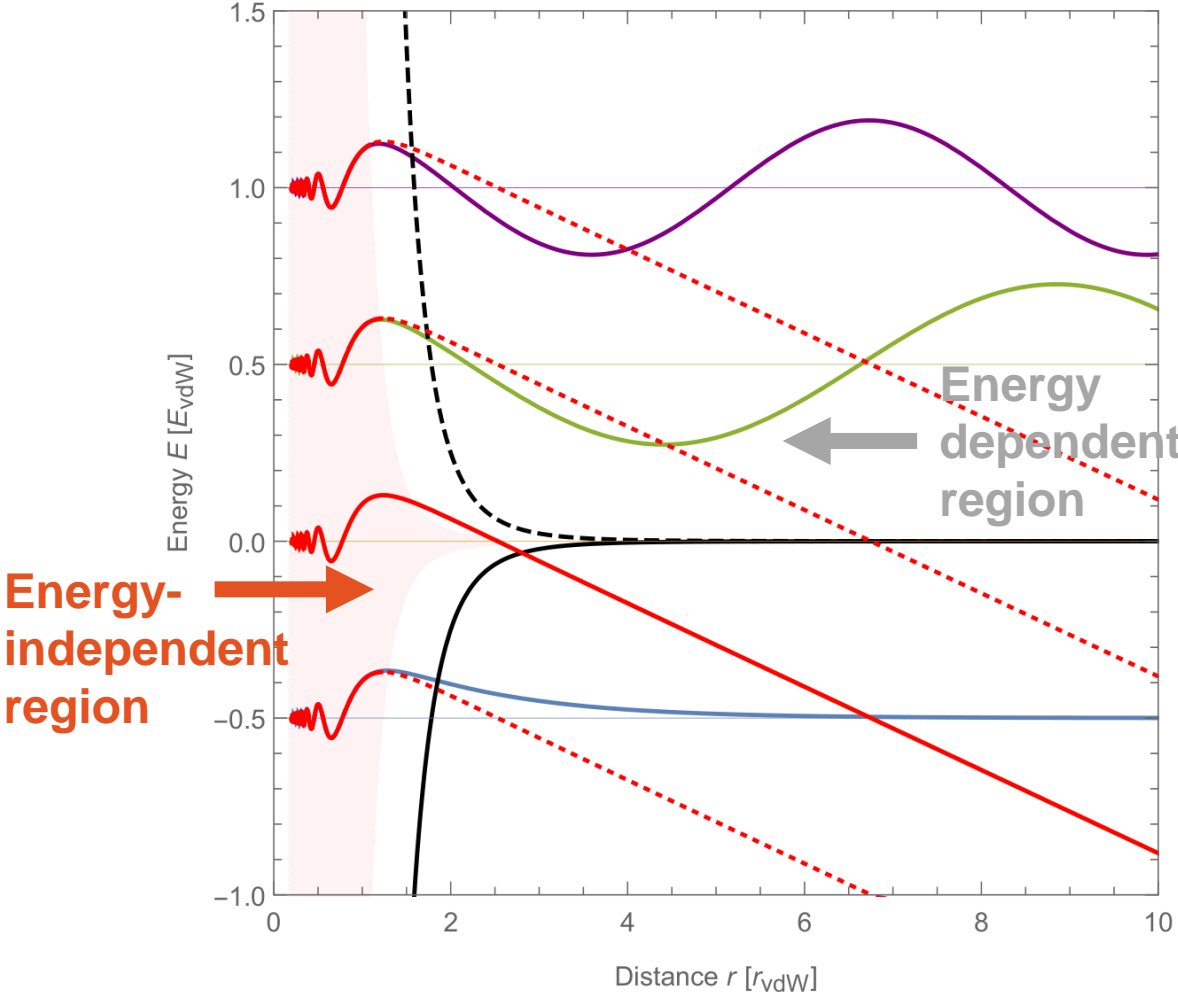
Answer: no.

For resonances described by a Quantum Defect Theory (QDT) such as

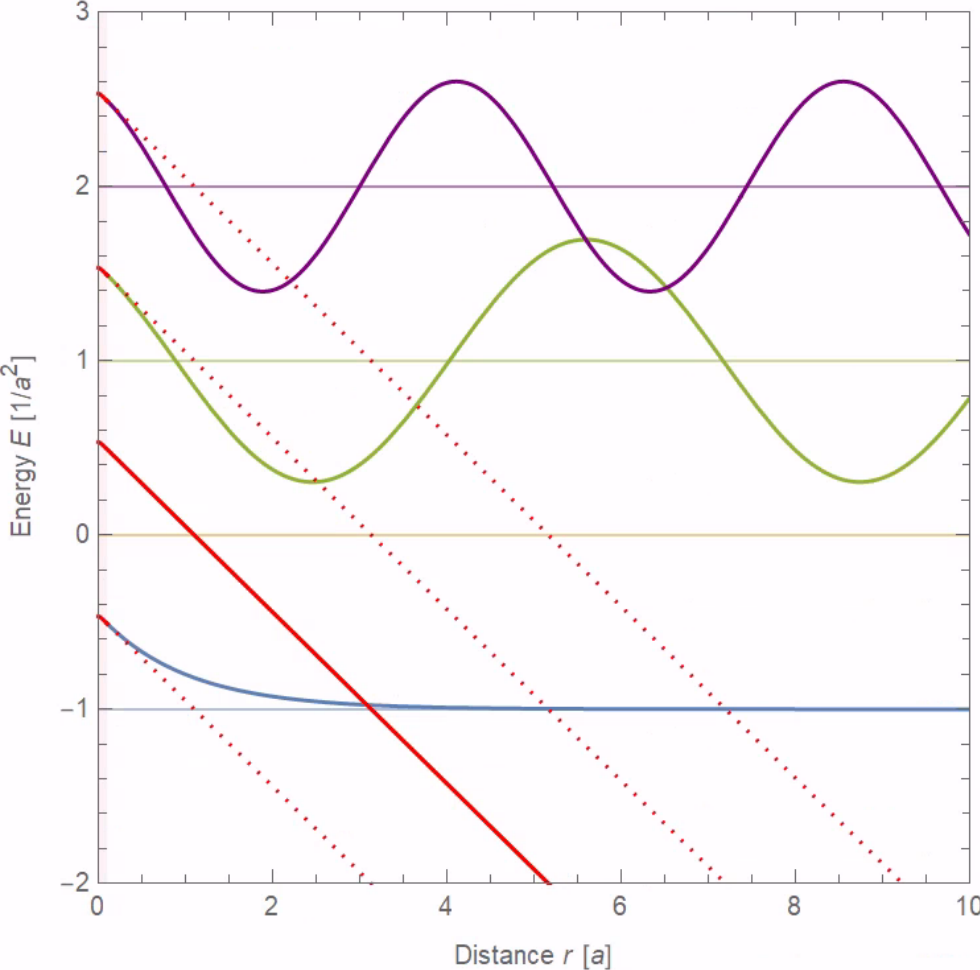
- All the atomic Feshbach resonances
- Hadron resonances very close to a threshold.

Quantum Defect Theory (QDT)

Van der Waals interaction



Short-range (Bethe-Peierls BC)



Quantum Defect Theory (QDT)

➔ The QDT gives explicit expressions for the shift and width:

$$\Delta(E) = \dots$$

$$\frac{\Gamma(E)}{2} = \dots$$



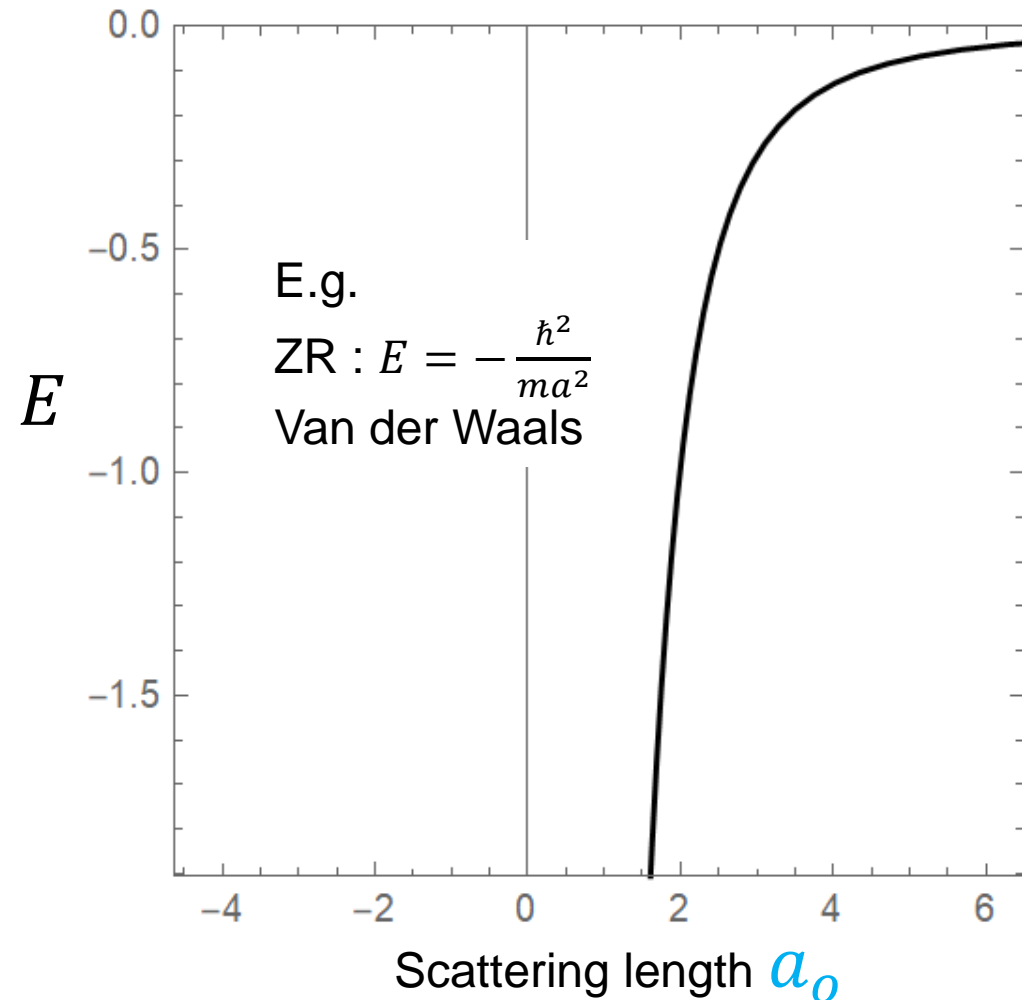
➔ The expressions should be **renormalised** in terms of observables

Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\text{shift}} \\ \Delta_0 + \frac{\gamma}{\lambda(E) - a_0}$$

dressed bare shift

Energy of an open-channel bound state

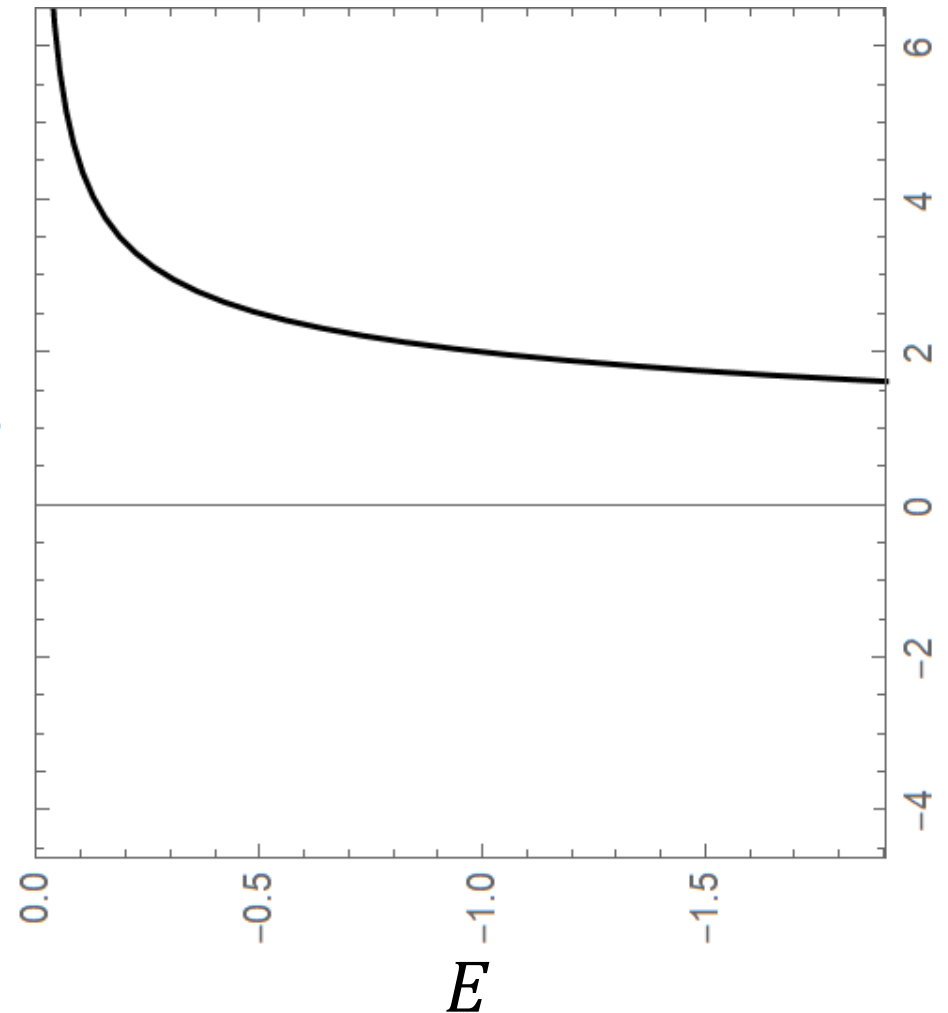


Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\Delta_0} + \frac{\gamma}{\lambda(E) - a_0}$$

Scattering length a_0

Energy of an open-channel bound state



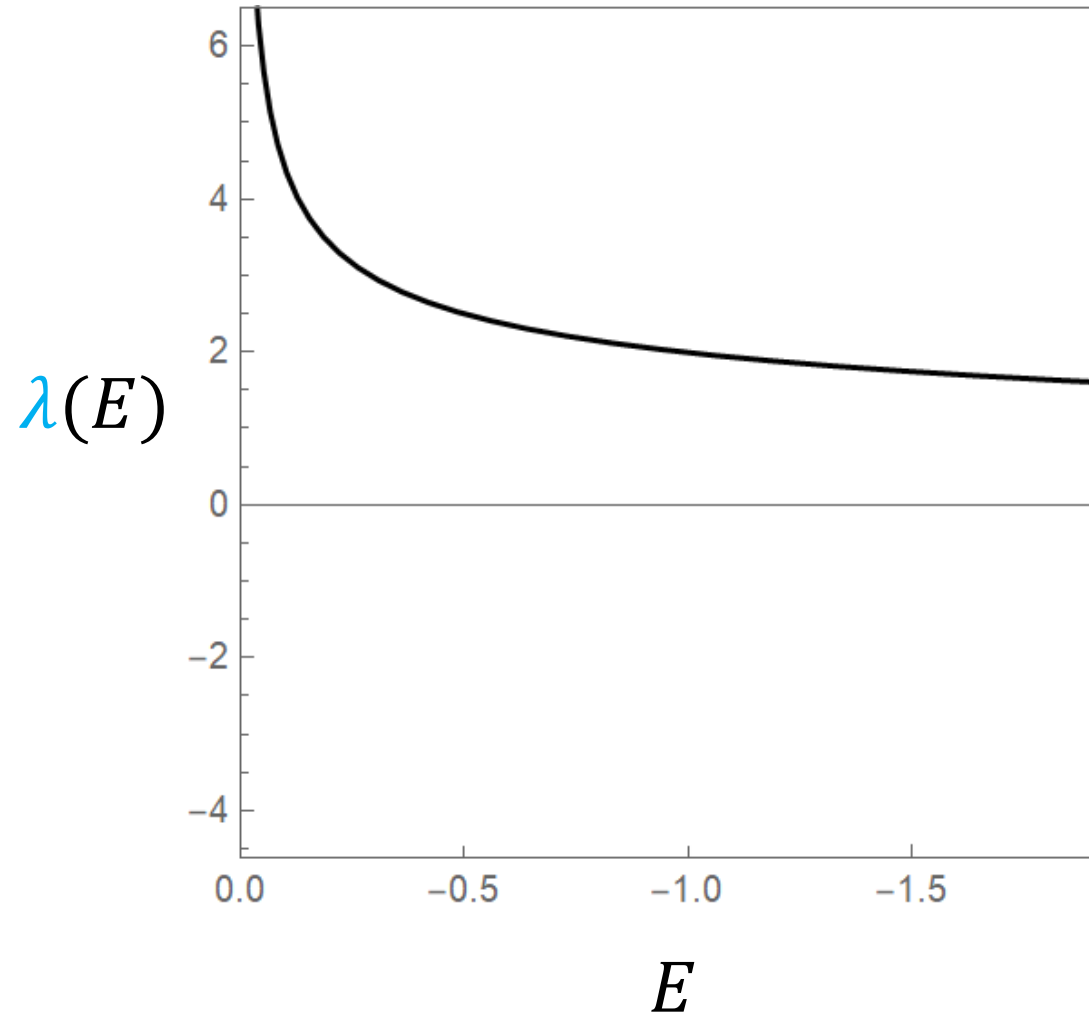
Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\Delta_0} + \frac{\gamma}{\lambda(E) - a_0}$$



$$E - \frac{\gamma}{\lambda(E) - a_0} = E_b + \Delta_0$$

Energy of an open-channel bound state



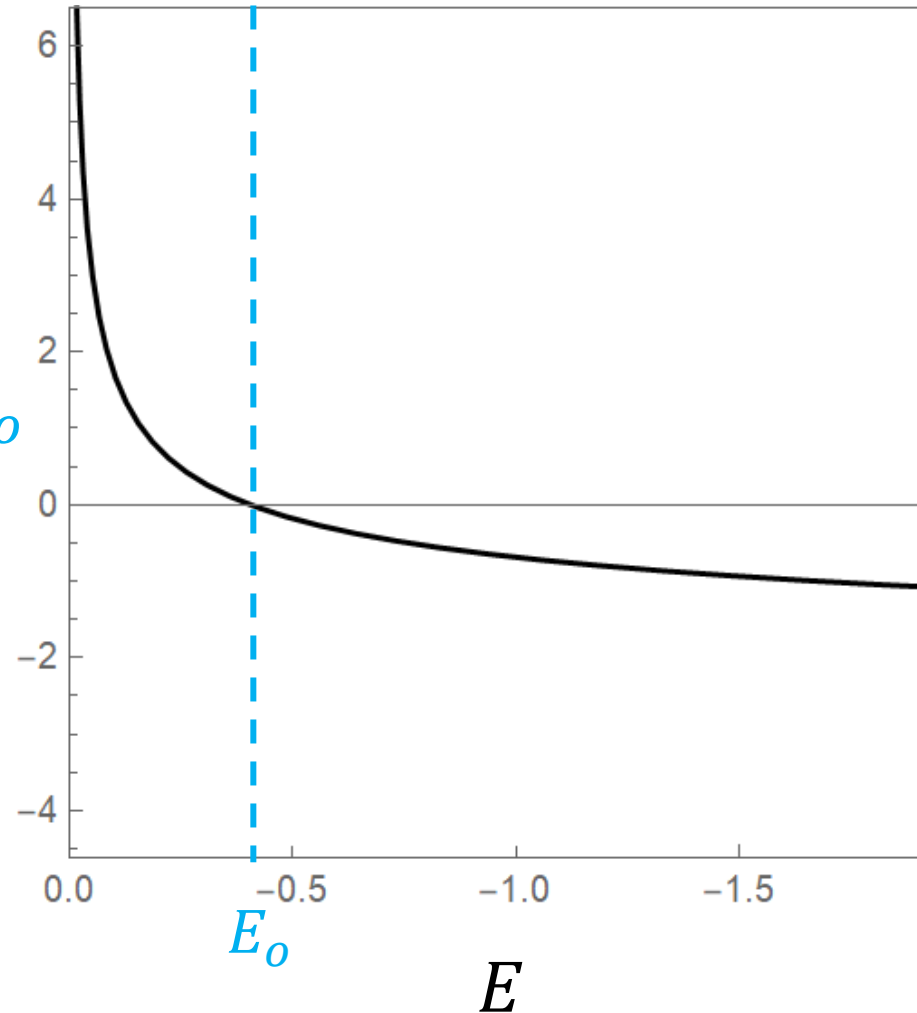
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$$\lambda(E) - a_0$$



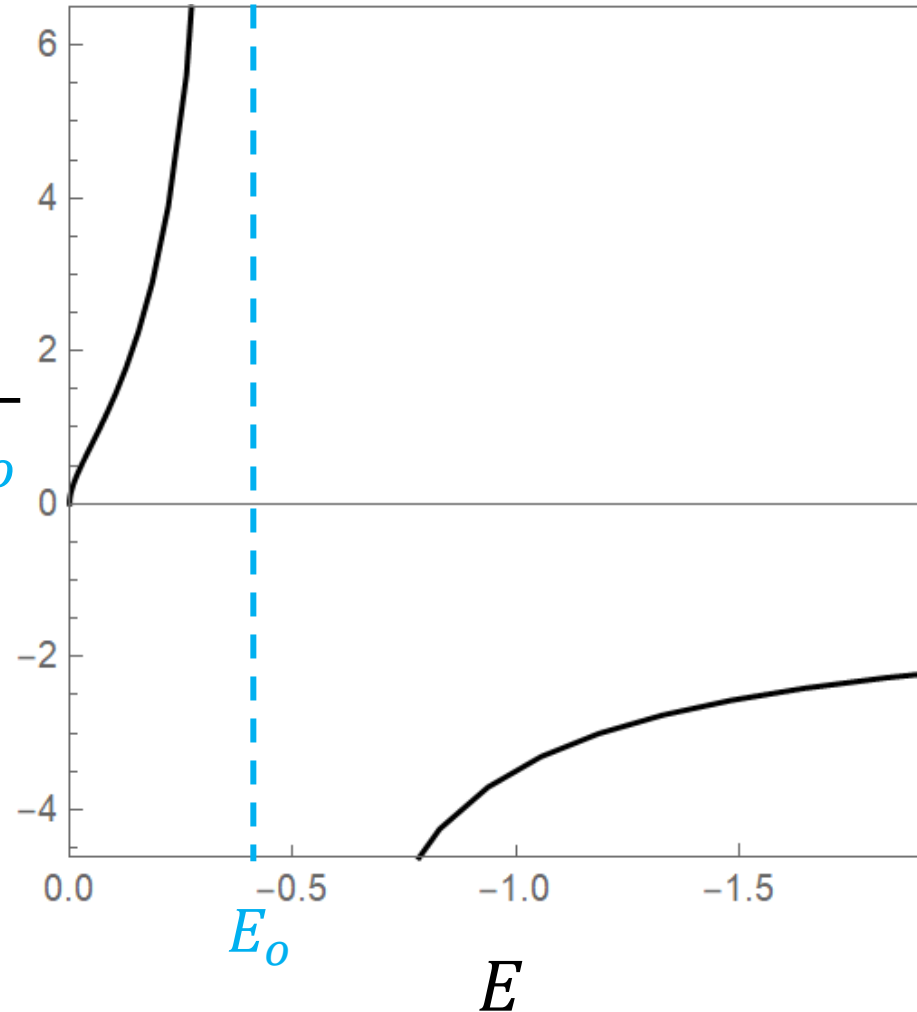
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$$\frac{\gamma}{\lambda(E) - a_0}$$

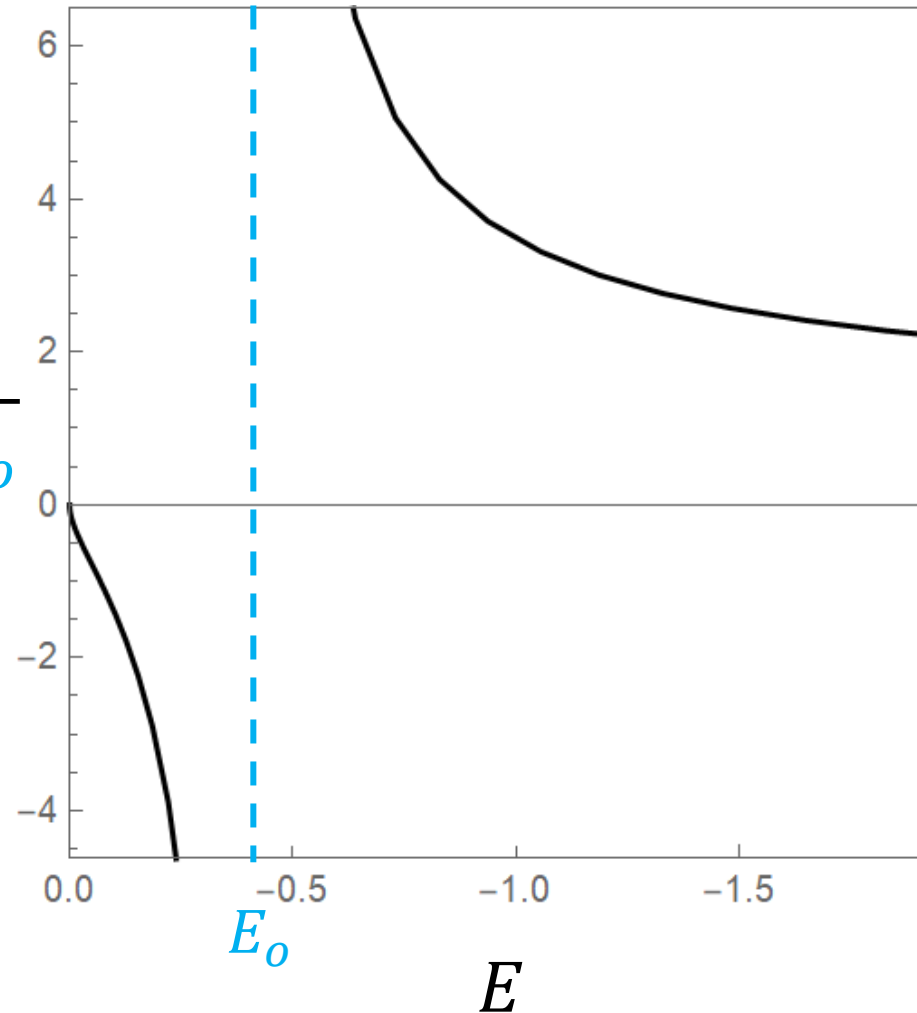


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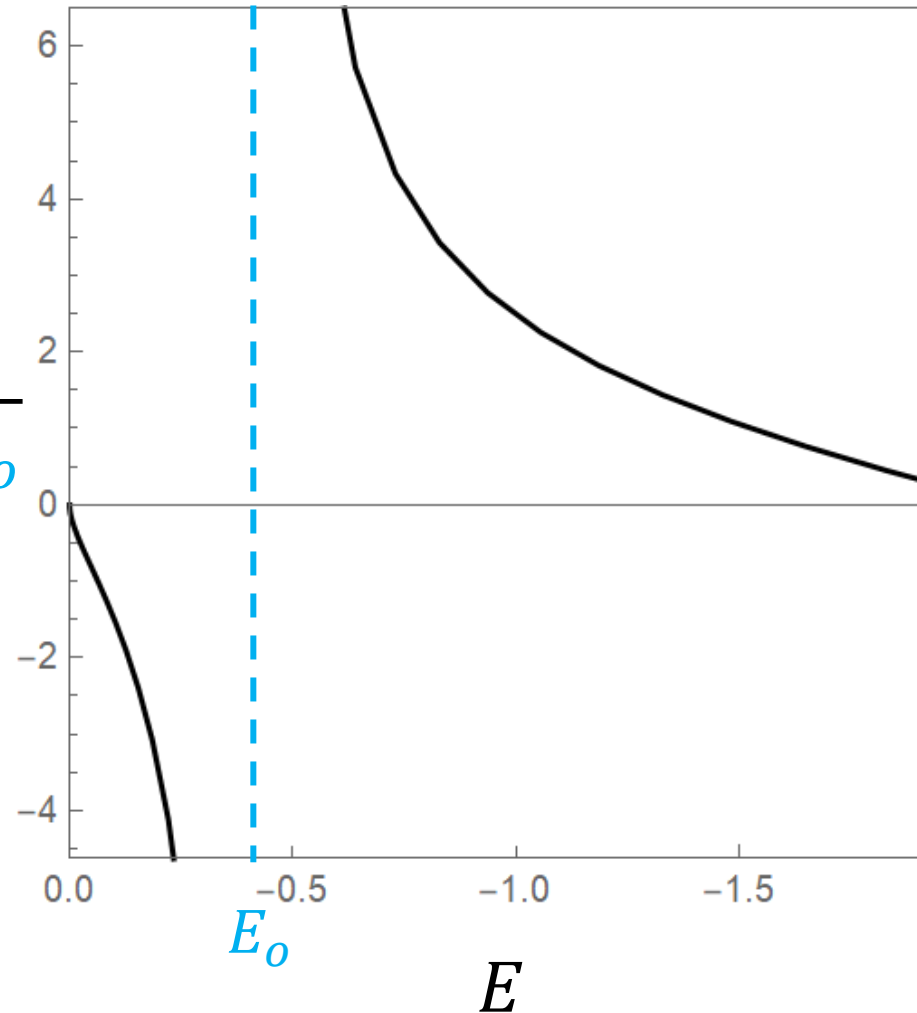
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$$E - \frac{\gamma}{\lambda(E) - a_0}$$



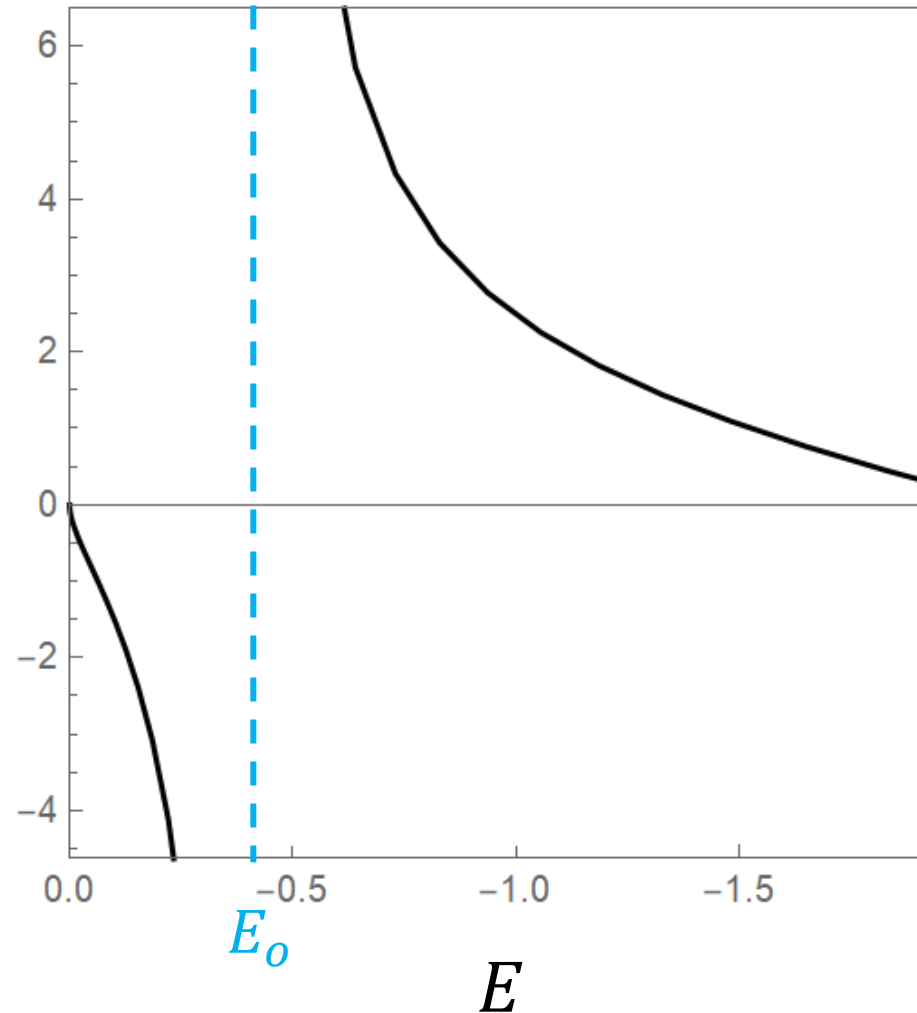
Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\Delta_0 + \frac{\gamma}{\lambda(E) - a_0}}$$



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$$E_b + \Delta_0$$

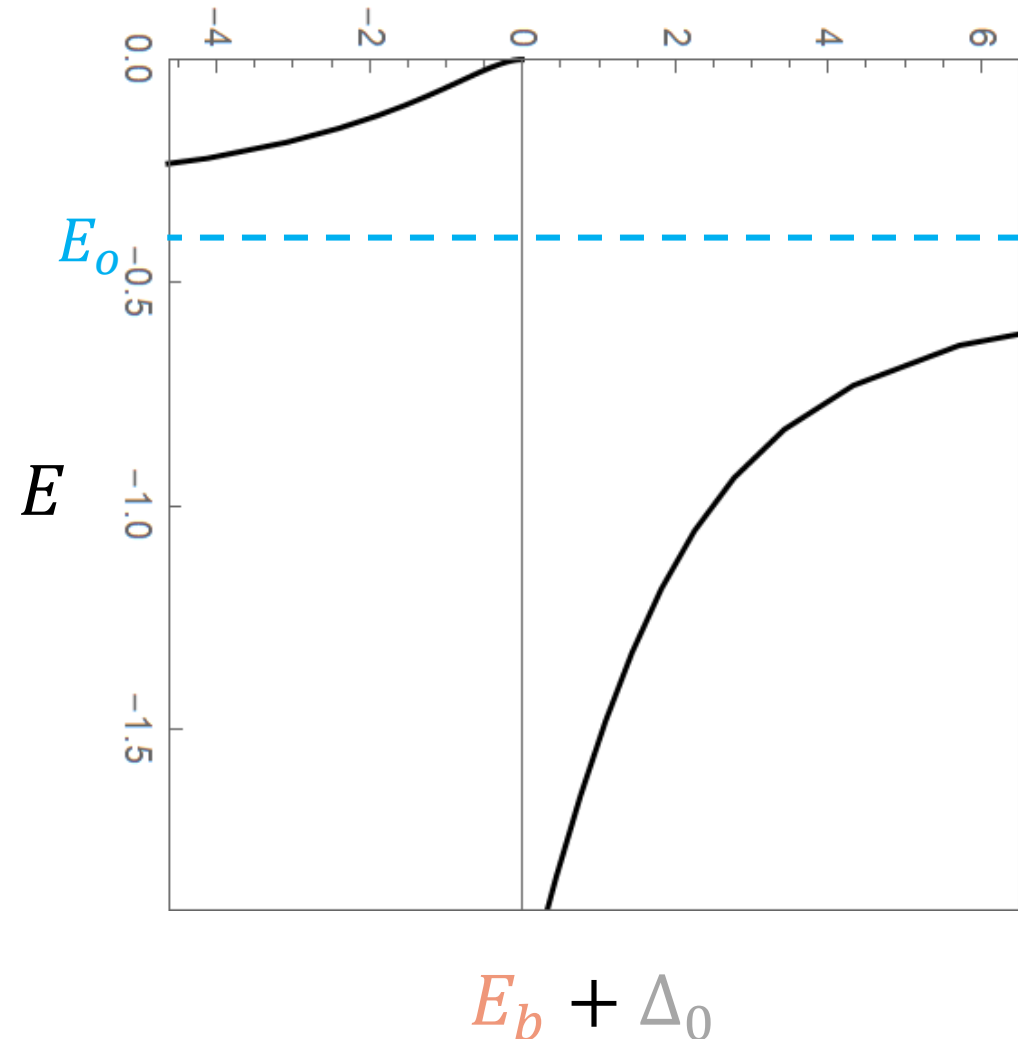


Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\Delta_0} + \frac{\gamma}{\lambda(E) - a_0}$$



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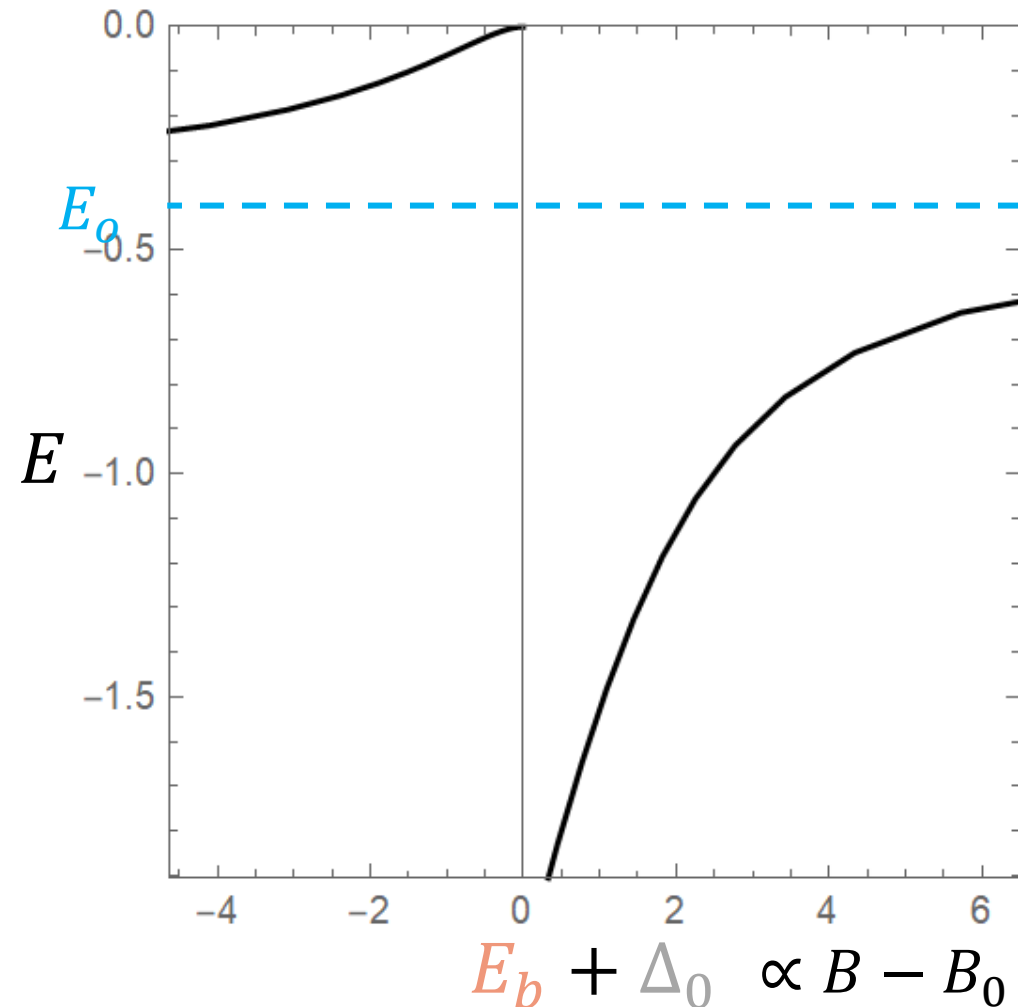


Renormalised Quantum Defect Theory

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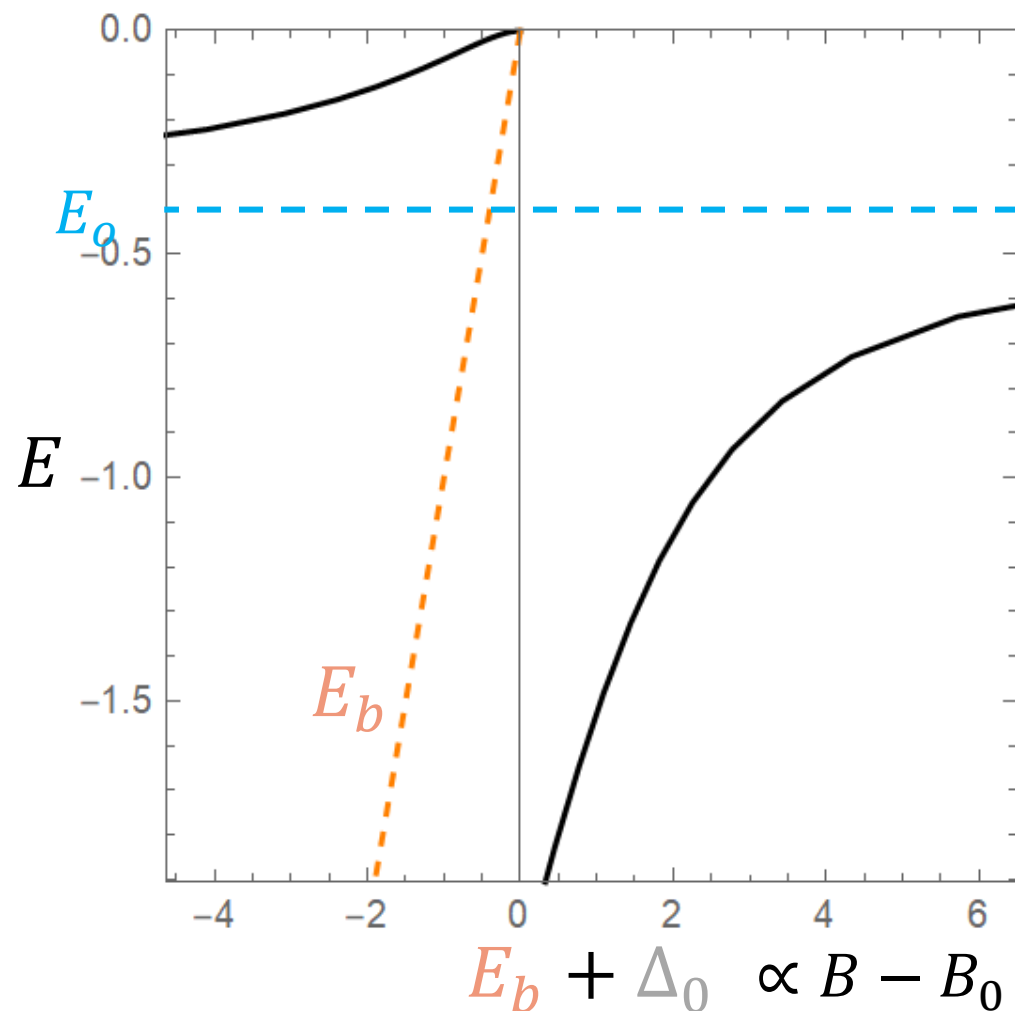


Renormalised Quantum Defect Theory

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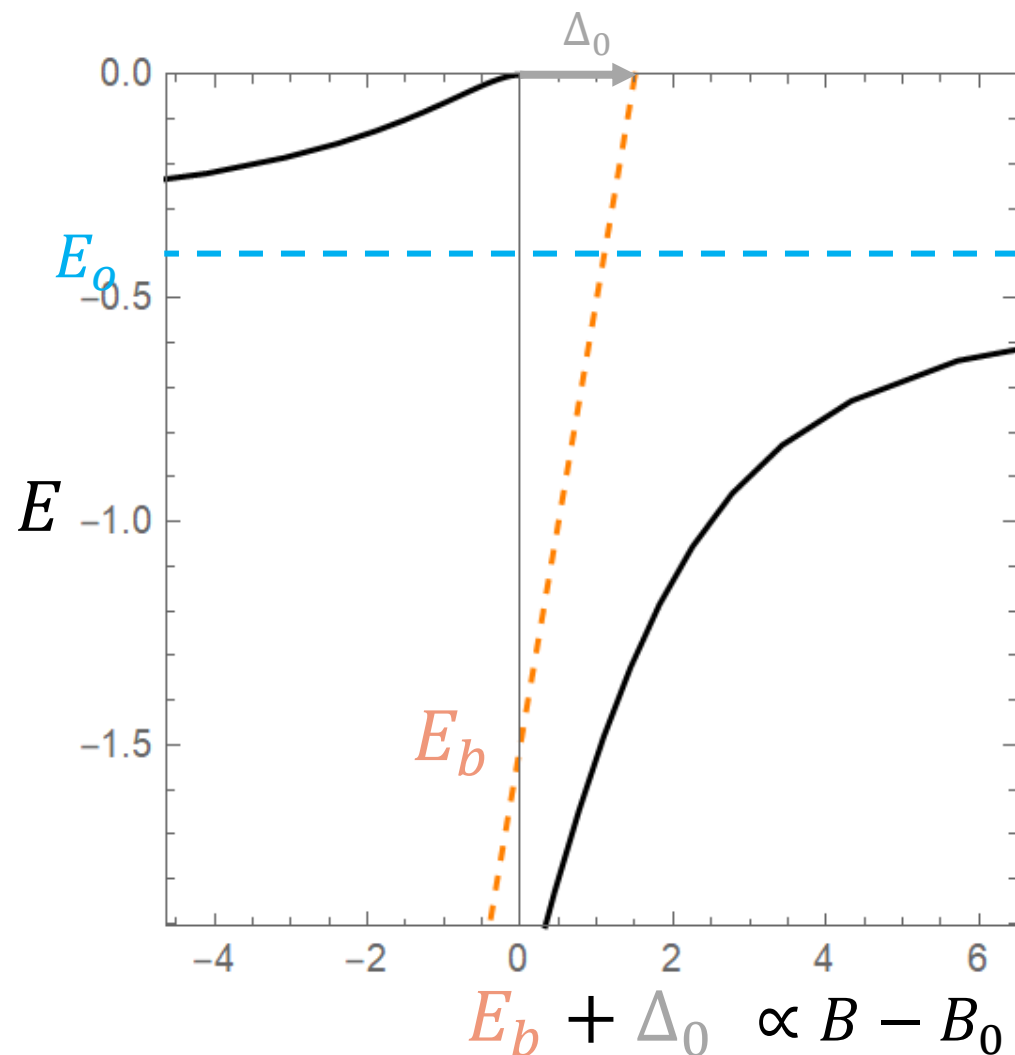


Renormalised Quantum Defect Theory

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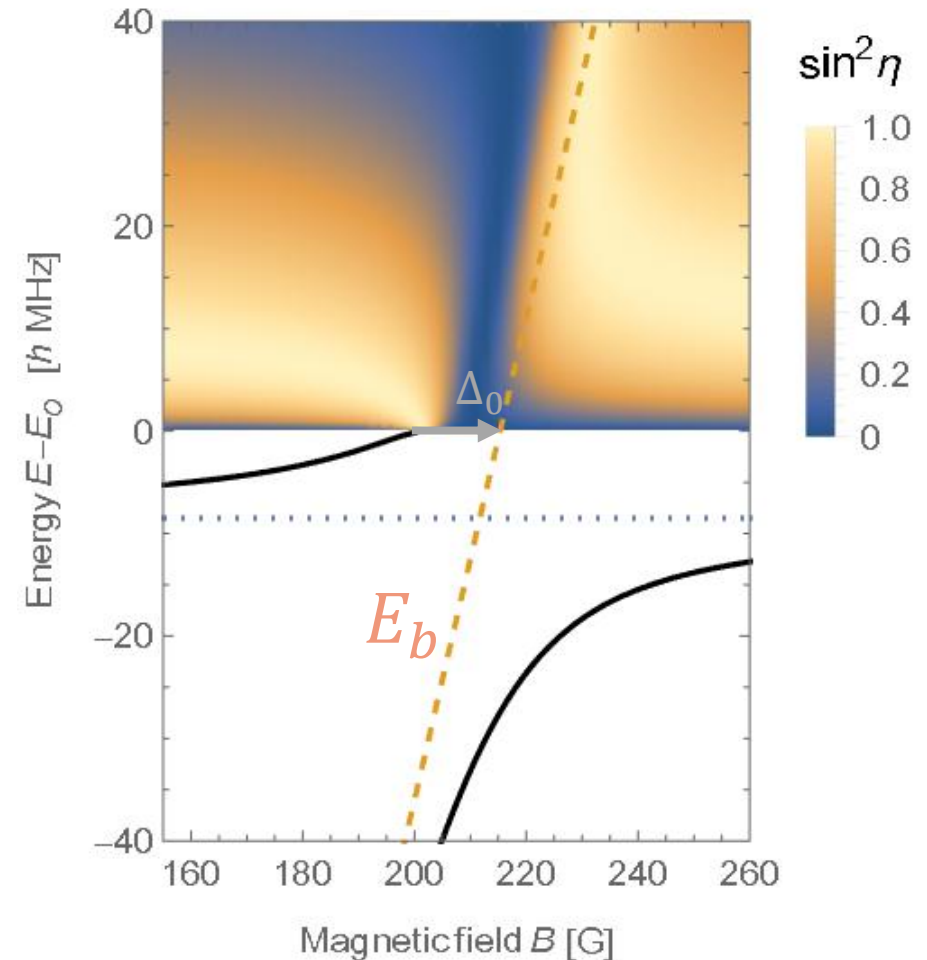


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Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta(E)}_{\Delta_0} + \frac{\gamma}{\lambda(E) - a_0}$$



Resonance of ^{40}K atoms near 202 G

Renormalised Quantum Defect Theory

$$E = E_b + \underbrace{\Delta_0 + \frac{\gamma}{\lambda(E) - a_o}}_{\text{shift}}$$

Expression for the shift Δ_0 :

~~$$\Delta_0 = \gamma \frac{a_o - \bar{a}}{\bar{a}^2 + (a_o - \bar{a})^2}$$~~

Chin, Grimm, Julienne, Tiesinga
Rev. Mod. Phys. 82, 1225 (2010)

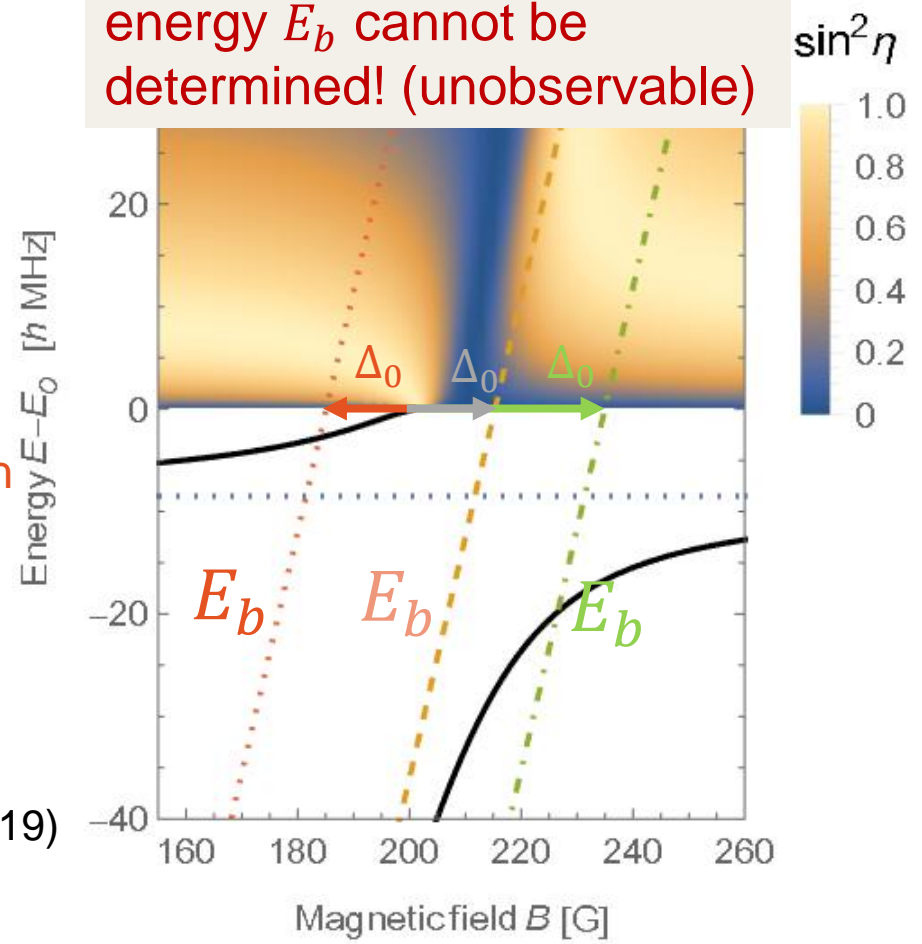
incorrect

Closed-channel scattering length

$$\Delta_0 = \frac{\gamma}{a_o - a_c}$$

Naidon & Pricoupenko
Phys. Rev. A 100, 042710 (2019)

The shift Δ_0 and the bare energy E_b cannot be determined! (unobservable)



Resonance of ^{40}K atoms near 202 G

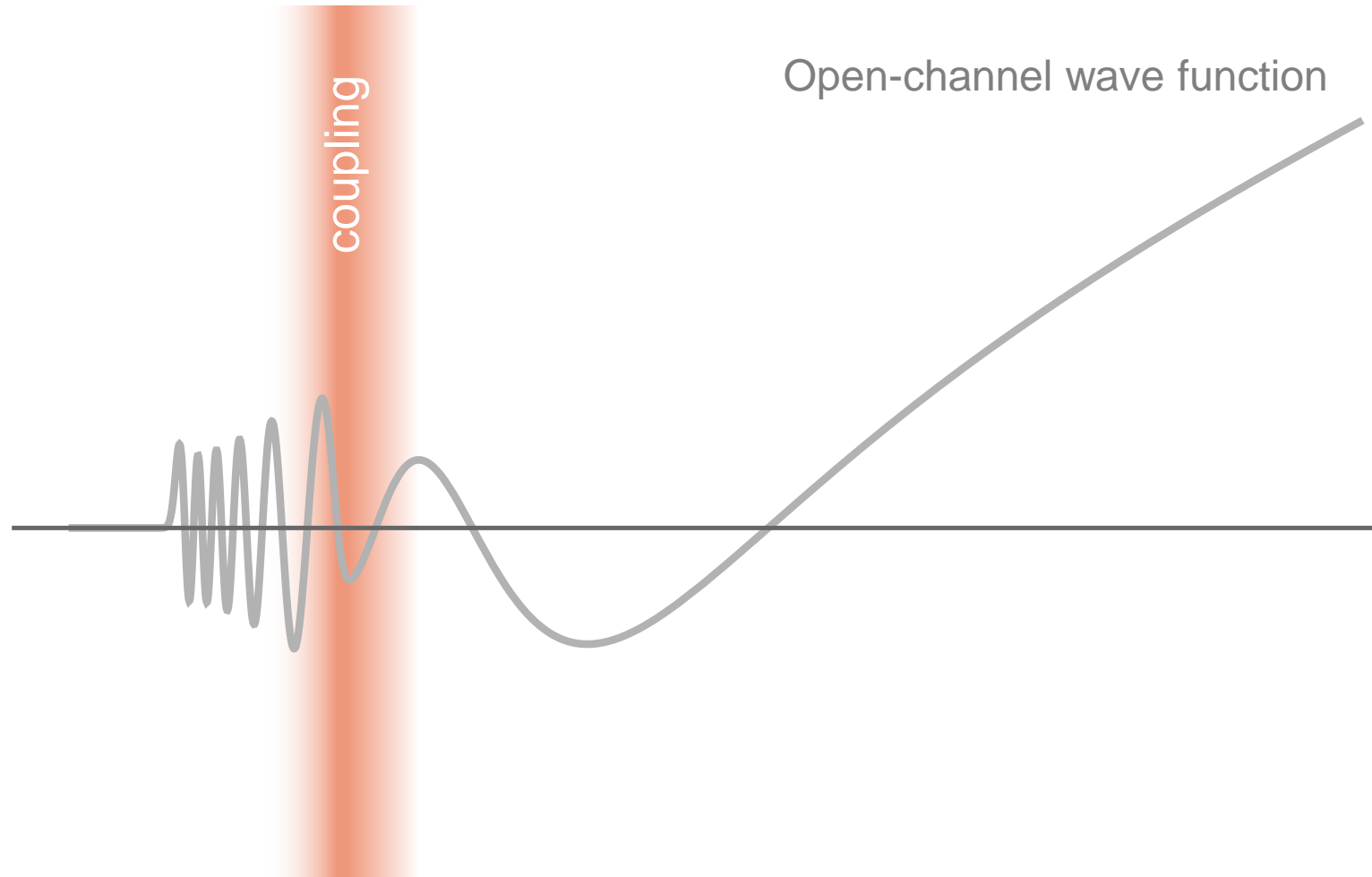
Preliminary conclusion

Closed-channel parameters (bare energy E_b , scattering length a_c) are **undetermined** by the **standard two-body observables** (scattering phase shifts, binding energy).

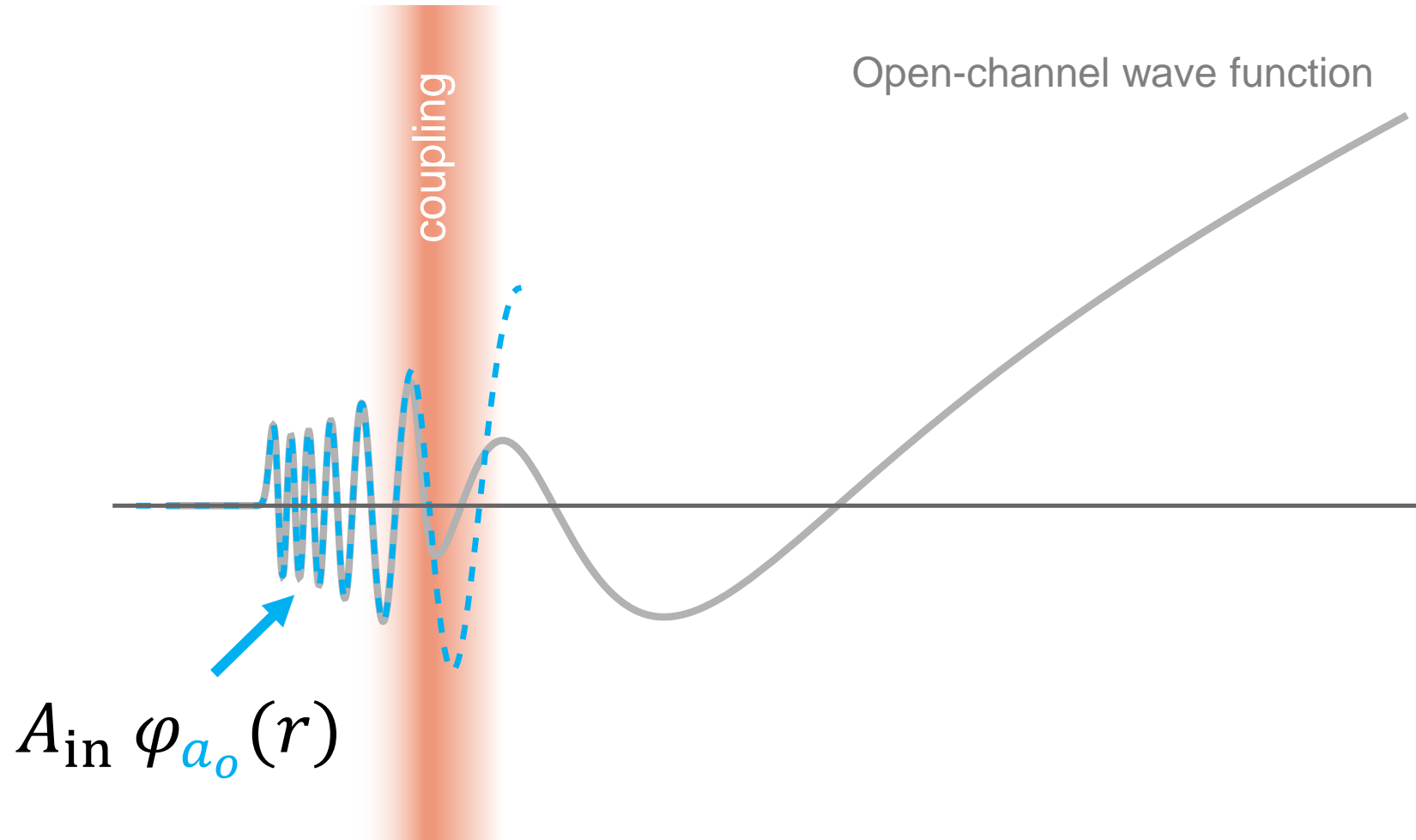
But one can also probe **short-range two-body observables** using a third particle :

- Three-body observables
- Photoassociation signal (probing with photon)

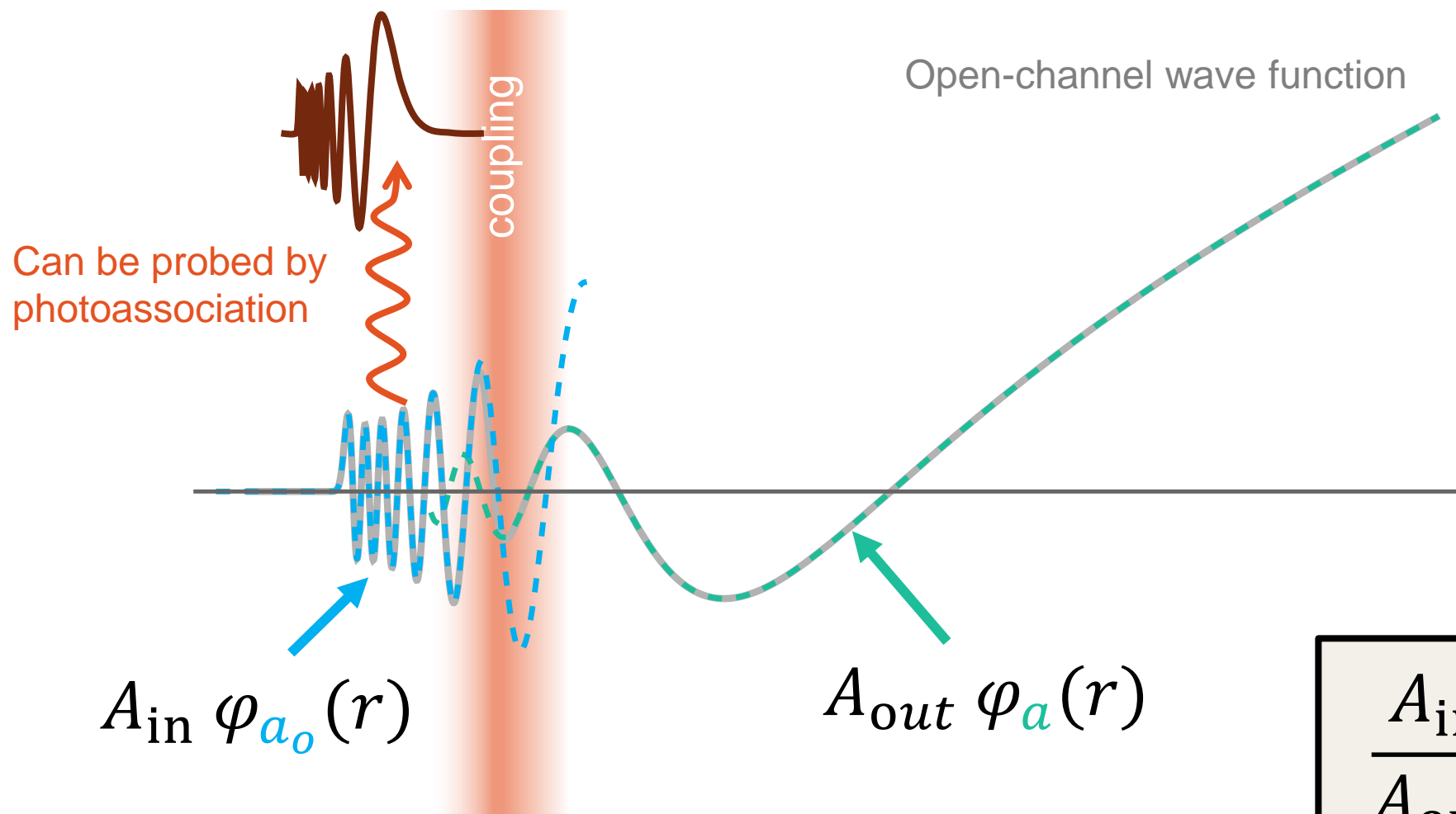
Short-range observables



Short-range observables



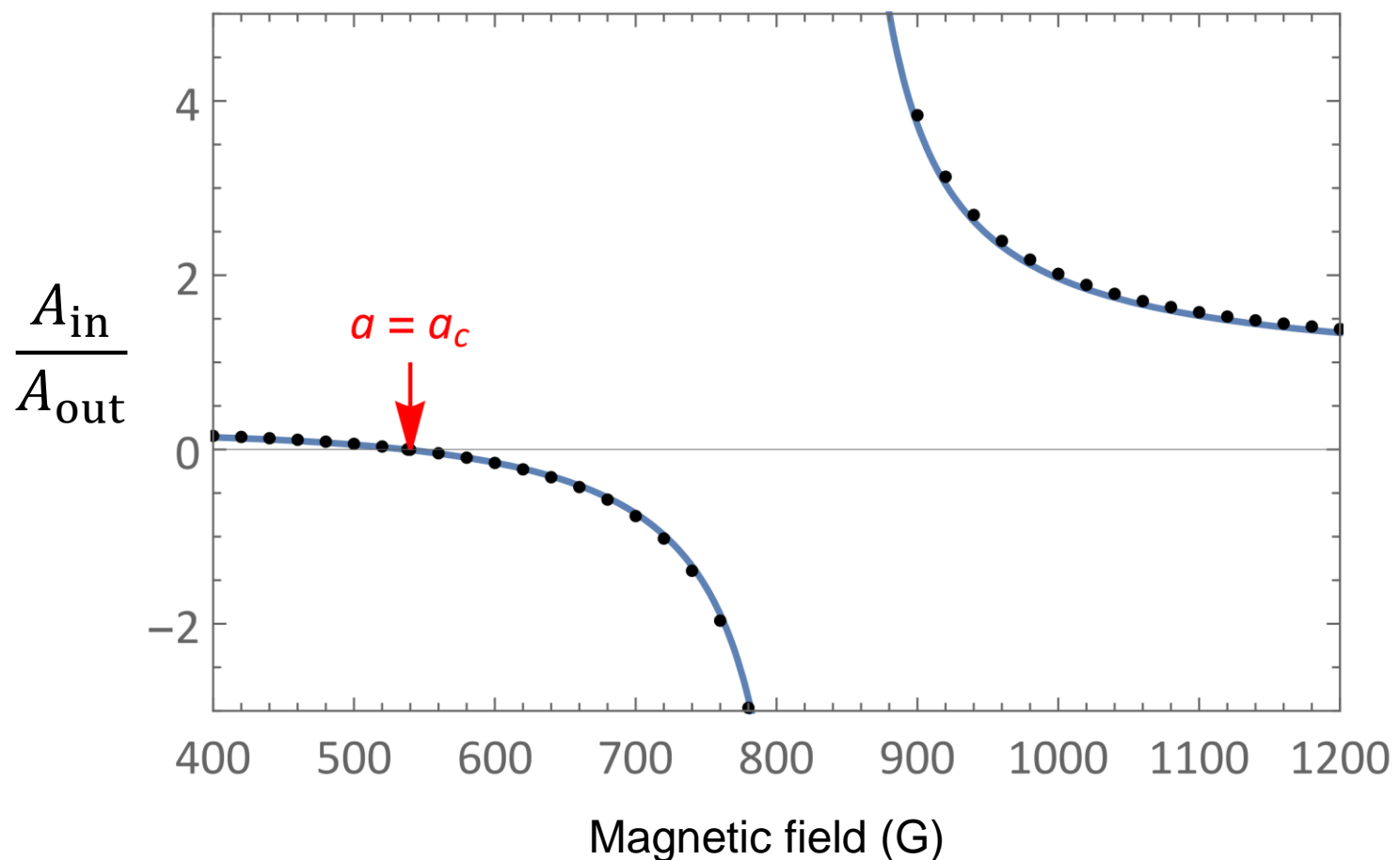
Short-range observables



$$\frac{A_{\text{in}}}{A_{\text{out}}} = \frac{a - a_c}{a_o - a_c}$$

Short-range observables

Resonance for ${}^6\text{Li}$ atoms near 834 G



● Coupled-channel calculation
(5 hyperfine channels)

— Formula $\frac{a - a_c}{a_o - a_c}$

$$\frac{A_{\text{in}}}{A_{\text{out}}} = \frac{a - a_c}{a_o - a_c}$$

Conclusion

The **renormalised quantum defect theory** gives a simple description of Feshbach resonances.

Standard two-body observables depend only the **open channel!**

Closed-channel parameters (bare energy E_b , scattering length a_c) can be revealed by **short-range two-body observables**. They may affect **three-body observables!**

arXiv:2403.14962